# Mathematical reflections on 

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I. The concept of locality revisited

## Locality principle

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Thus, one can separate events located in different regions of space-time and should be able to measure them independently.

- Propose a mathematical framework which encompasses the main features of the locality principle in QFT;
- use this framework to carry out renormalisation in accordance with the locality principle.


## Causal separation

## Light cone, past and future

In the Minkowski space $\left(\mathbb{R}^{d}, g\right)$, where $g(x, y)=-x_{0} y_{0}+\sum_{j=1}^{d-1} x_{j} y_{j}$ is the Lorentzian scalar product,

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(picture downloaded from Wikipedia)

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(picture downloaded from Wikipedia)

Two sets $S_{1}$ and $S_{2}$ are causally separated $\left(S_{1} \| S_{2}\right)$ if and only if $S_{i}$ does not lie in the future of $S_{j}$ for $i \neq j$.

## Locality in axiomatic QFT

The Wightman field $\varphi: \mathcal{S}\left(\mathbb{R}^{d}\right) \rightarrow \mathcal{O}(H)$ obeys the locality axiom
$\operatorname{Supp}\left(f_{1}\right) \| \operatorname{Supp}\left(f_{2}\right) \Longrightarrow\left[\varphi\left(f_{1}\right), \varphi\left(f_{2}\right)\right]=0$.

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The (relative) scattering matrix $S_{f}$ satisfies the locality condition
$\operatorname{Supp}\left(f_{1}\right) \| \operatorname{Supp}\left(f_{2}\right) \Longrightarrow S_{f}\left(f_{1}+f_{2}\right)=S_{f}\left(f_{1}\right) S_{f}\left(f_{2}\right)$

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& \Longrightarrow\left[S_{f}\left(f_{1}\right), S_{f}\left(f_{2}\right)\right]=0 . \tag{2}
\end{align*}
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## Mathematical interpretation

We introduce two binary relations

- on sets:

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O_{1} \top^{\prime} O_{2}: \Leftrightarrow\left[O_{1}, O_{2}\right]=0, \tag{3}
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f_{1} \top f_{2} \Longrightarrow S_{f}\left(f_{1}+f_{2}\right)=S_{f}\left(f_{1}\right) S_{f}\left(f_{2}\right) . \tag{6}
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# II. Locality as a symmetric binary relation 

## Definition of locality

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## (almost-)Separation of supports

Let $U \subset \mathbb{R}^{n}$ be an open subset and $\epsilon \geq 0$. Two functions $\phi, \psi$ in $\mathcal{D}(U)$ are independent i.e., $\phi \top_{\epsilon} \psi$ whenever

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## Further examples

Probability theory: independence of events
Given a probability space $\mathcal{P}:=(\Omega, \Sigma, P)$ and two events $A, B \in \Sigma$ :
$A \top B \Longleftrightarrow P(A \cap B)=P(A) P(B)$.

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## Geometry: transversal manifolds

Given two submanifolds $L_{1}$ and $L_{2}$ of a manifold $M$ :

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L_{1} \top L_{2}: \Longleftrightarrow L_{1} \pitchfork L_{2} \Longleftrightarrow T_{x} L_{1}+T_{x} L_{2}=T_{x} M \quad \forall x \in L_{1} \cap L_{2} .
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## Number theory: coprime numbers

Given two positive integers $m, n$ in $\mathbb{N}$ :

$$
m 丁 n \Longleftrightarrow m \wedge n=1 .
$$

## category

## Locality structures

- set $X \rightsquigarrow$ locality set $(X, \top)$; the polar set of $U$ is $U^{\top}:=\{x \in X, x \top u \quad \forall u \in U\}$
- semi-group $\left(G, m_{G}\right) \rightsquigarrow$ locality semi-group $\left(G, m_{G}, \top\right)$ ( $U \subset G \Longrightarrow U^{\top}$ semi-group);


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- algebra $\left(A,+, \cdot, m_{A}\right) \rightsquigarrow$ locality algebra $\left(A,+, \cdot, m_{A}, \top\right)$.


## Locality morphisms: $f:\left(X, \top_{X}\right) \rightarrow(Y, \top Y)$

- locality map:
$(f \times f)\left(\top_{x}\right) \subset \top_{y} \quad$ or equivalently $\quad x_{1} \top_{x} x_{2} \Longrightarrow f\left(x_{1}\right) \top_{y} f\left(x_{2}\right)$;


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- locality semi-group morphism $f:\left(X, m_{X}, \top_{X}\right) \rightarrow\left(Y, m_{Y}, \top_{Y}\right)$ : $f$ is a locality map and $x_{1} \top x x_{2} \Longrightarrow f\left(m_{X}\left(x_{1}, x_{2}\right)\right)=m_{Y}\left(f\left(x_{1}\right), f\left(x_{2}\right)\right)$ etc...
III. Locality relations are ubiquitious


## Local functionals

are functionals $F$ on test functions (fields) $\varphi$ of the form $F(\varphi)=\int_{M} f\left(j_{x}^{k}(\varphi)\right) d x$ (here $j_{x}^{k}(\phi)$ is the $k$-th jet of $\phi$ at $x$ ): The localised version at $\varphi$ :

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\begin{equation*}
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Hammerstein property/partial additivity simiar to a causality condition on S -matrices of [Epstein, Glaser (1973)], [Bogoliubov, Shirkov (1959))], [Stückelberg (1950, 1951)]

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\begin{equation*}
\varphi_{1} \top n \varphi_{2} \Longrightarrow F\left(\varphi_{1}+\varphi+\varphi_{2}\right)=F\left(\varphi_{1}+\varphi\right)-F(\varphi)+F\left(\varphi+\varphi_{2}\right) \quad \forall \varphi . \tag{8}
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Comparing the two [Brouder, Dang, Laurent-Gengoux, Rejzner (2018)] Provided $D_{\varphi} F$ can be represented as a function $\nabla_{\varphi} F$ such that the map $\varphi \mapsto \nabla_{\varphi} F$ is smooth, then $\quad(8) \Longleftrightarrow(7)$.

## and singularities

## Separation of wavefront sets

We define two locality relations on on $\mathcal{D}^{\prime}(U), U \subset \mathbb{R}^{n}$ :

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where we have set $\mathrm{WF}^{\prime}(v):=\left\{(x,-\xi) \in U \times\left(\mathbb{R}^{n} \backslash\{0\}\right) \mid(x, \xi) \in \mathrm{WF}(v)\right\}$.

Counterexample
Distributions can be independent for $T^{\mathrm{WF}}$ and not for $T^{\text {sing }}$. We have $v_{1} T^{\text {sing }} v_{2} \Longrightarrow v_{1} \top^{W F} v_{2}$

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Distributions can be independent for $T^{W F}$ and not for $T^{\text {sing }}$. We have $v_{1} \top^{\text {sing }} v_{2} \Longrightarrow v_{1} \top^{W F} v_{2}$ but not conversely. The wavefront sets of $\nu_{\mathbf{1}}(\phi):=\int_{\mathbb{R}^{2}} \phi(0, y) d y$ and $\nu_{\mathbf{2}}(\phi):=\int_{\mathbb{R}^{2}} \phi(x, 0) d x$ read $\mathrm{WF}\left(\nu_{\mathbf{1}}\right)=\{((0, y) ;(\lambda, 0)) \mid y \in \mathbb{R}, \lambda \in \mathbb{R} \backslash\{0\}\} \quad ; \quad \mathrm{WF}\left(\nu_{\mathbf{2}}\right)=\{((x, 0) ;(0, \mu)) \mid x \in \mathbb{R}, \mu \in \mathbb{R} \backslash\{0\}\}$, so $\nu_{1} \top^{\mathrm{WF}} \nu_{2}$ but $\nu_{1} \top^{\text {sipg }} \nu_{2}$.

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## Partial product of pseudodifferential operators of non-integer order

We equip $\Psi_{\text {pgh }}^{\notin \mathbb{Z}}$ (the canonical trace $T \mathrm{R}$ is well defined) with the locality relation $A_{1} \top \notin \mathbb{Z} A_{2}: \Leftrightarrow\left(\operatorname{ord}\left(A_{1}\right)+\operatorname{ord}\left(A_{2}\right) \notin \mathbb{Z}\right) \Rightarrow\left(\mathrm{TR}\left(\left[A_{1}, A_{2}\right]\right)=0\right)$.

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(1) When is the quotient $V / W$ of a locality vector space $(V, \top)$ by a linear subspace $W$, a locality vector space if equipped with the quotient locality relation $T$ given by the final locality relation:

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IV. Evaluating meromorphic germs at poles in QFT

## Functions of several variables in QFT

## Speer's analytic renormalisation [JMP 1967] revisited

Eugene Speer considers Feynman amplitudes given by the coefficients of the perturbation-series expansion of the $S$ matrix in a Lagrangian field theory (with non zero mass).

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## Excerpt of Speer's article

In this paper we apply a method of defining divergent quantities which was originated by Riesz and has been used in various contexts by many authors. [....] We find it necessary to consider functions of several complex variables $z_{1}, \cdots, z_{k}$, one associated with each line of the Feynman graph. The main difficulty is the extension of the above [Riesz's] treatment of poles to the more complicated singularities which occur in several complex variables...

## Brain teaser

(We assume the poles are at zero)
Speer shows [Theorem 1] that the divergent expressions lie in the filtered algebra $\mathcal{M}^{\text {Feyn }}\left(\mathbb{C}^{\infty}\right):=\cup_{k=1}^{\infty} \mathcal{M}^{\text {Feyn }}\left(\mathbb{C}^{k}\right)$ consisting of Feynman functions $f: \mathbb{C}^{k} \rightarrow \mathbb{C}$,

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## Evaluating a fraction with a linear pole at zero

$$
f\left(z_{1}, z_{2}\right)=\left.\frac{z_{1}-z_{2}}{z_{1}+z_{2}}\right|_{z_{1}=0, z_{2}=0}=\left\{\begin{array}{c}
1 \text { or }-1 ? \\
0 ? \\
10000 ?
\end{array}\right.
$$

## V. Locality on meromorphic germs comes to the rescue

## Locality on multiparameter meromorphic germs

Multiparameter meromorphic germs with linear poles

- $\mathcal{M}\left(\mathbb{C}^{k}\right) \ni f=\frac{h\left(\ell_{1}, \cdots, \ell_{n}\right)}{L_{1}^{1} \cdots L_{n}^{n}}, h$ holomorphic germ, $s_{i} \in \mathbb{Z}_{\geq 0}$,


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## Locality: separation of variables

On $\mathcal{M}\left(\mathbb{C}^{\infty}\right)=\bigcup_{k \in \mathbb{N}} \mathcal{M}\left(\mathbb{C}^{k}\right), f_{1} Q^{\top} f_{2} \Longleftrightarrow \operatorname{Dep}\left(f_{1}\right) \perp \operatorname{Dep}\left(f_{2}\right)$.

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## Back to the locality principle in QFT

We consider $\mathcal{M}:=\mathcal{M}\left(\mathbb{C}^{\infty}\right):=\cup_{k=1}^{\infty} \mathcal{M}\left(\mathbb{C}^{k}\right)$ consisting of meromorphic functions/germs $f: \mathbb{C}^{k} \rightarrow \mathbb{C}$ with linear poles at zero,

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## Principle of locality revisited: locality evaluators

$f \perp^{Q} g \Longrightarrow \mathcal{E}(f \cdot g)=\mathcal{E}(f) \mathcal{E}(g)$ for two meromorphic germs $f$ and $g$ in an appropriate subalgebra $\mathcal{M}^{\bullet}$ of $\mathcal{M}$.

## Speer's generalised evaluators

Reminder: Meromorphic germs in $\mathcal{M}^{\text {Feyn }}\left(\mathbb{C}^{k}\right)$ have linear poles $L_{i}=\sum_{j_{i} \in J_{i}} j_{i}$.
Speer's evaluators consist of a family $\mathcal{E}=\left\{\mathcal{E}_{k}, \in \mathbb{N}\right\}$ of linear forms $\mathcal{E}_{k}: \mathcal{M}^{\mathrm{Feyn}}\left(\mathbb{C}^{k}\right) \rightarrow \mathbb{C}$, compatible with the filtration, which fulfill the following conditions

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(3) $\mathcal{E}$ is invariant under permutations of the variables $\mathcal{E}_{k} \circ \sigma^{*}=\mathcal{E}_{k}$ for any $\sigma \in \Sigma_{k}$, with $\sigma^{*} f\left(z_{1}, \cdots, z_{k}\right):=f\left(z_{\sigma(1)}, \cdots, z_{\sigma(k)}\right)$;

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(1) (extend $\left.\mathrm{ev}_{0}\right) \mathcal{E}$ is the ordinary evaluation $\mathrm{ev}_{0}$ at zero on holom. germs;
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## Speer's generalised evaluators

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Drawback: Speer's approach depends on the choice of coordinates

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## Orthogonal projection

$\perp^{Q}$ induces a splitting [Berline and Vergne 2005, Guo, Zhang, S.P. 2015]

$$
\mathcal{M}^{\bullet}=\mathcal{M}_{+} \oplus^{Q} \mathcal{M}_{-}^{\bullet} \quad \text { and } \quad \pi_{+}{ }^{Q}: \mathcal{M}^{\bullet} \longrightarrow \mathcal{M}_{+}
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## VI. Classification of locality evaluators

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- a locality isomorphism $u \mapsto x_{u}$ between the locality algebra generated by Chen-type poles $L_{i}=\sum_{j=1}^{i} \ell_{u_{j}}=\ell_{u_{1}}+\cdots \ell_{u_{i}}$ with $u \top v \Longrightarrow \ell_{u} \perp^{Q} \ell_{v}$ and a certain locality shuffle algebra.
- Conclusion: $\mathcal{M}^{\text {Chen }}\left(\mathcal{M}^{\mathrm{Feyn}}\right)$ are locality polynomial algebras with locality "Lyndon fractions" as locality generators.


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Since $\mathcal{M}^{\text {Chen }}$, resp. $\mathcal{M}^{\text {Feyn }}$ are $\perp$-local polynomial algebras, a generalised evaluator is uniquely determined by its values on the free generators.

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If $\mathcal{M}^{\bullet}$ is a free polynomial locality-algebra generated by $\mathcal{S}^{\bullet}$, then $T \in \operatorname{Gal}^{-}\left(\mathcal{M}^{\bullet} / \mathcal{M}_{+}\right)$ is uniquely determined by $\left\{T(S), S \in \mathcal{S}^{\bullet}\right\}$ :
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THANK YOU FOR YOUR ATTENTION!
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