

# Finite controllability of $\infty$ -d quantum systems: Approximate controllability of trapped-ion quantum states



Chitra Rangan

Department of Physics, University of Windsor

# Outline

- Work done in collaboration with Tony Bloch (Michigan) and Roger Brockett (Harvard)
- Finite controllability
- E.g.: spin-half coupled to QHO
  - Can show approx. controllability
- Two similar systems that are not finite controllable

# Finite controllability

## Definition

**Given**

**-a system, and**

**-a nested set of finite dimensional subspaces**

**it will be said to be finitely controllable if**

**- it can be transferred from any point in one of the subspaces to any other point in that subspace**

**- with a trajectory lying entirely within the subspace.**

# Finite controllability theorem

**Consider a complex Hilbert space  $X$  together with a nested set of finite-dimensional subsets**

$$H = \{H_1 \subset H_2 \subset H_3 \dots\}$$

**Consider** 
$$i\dot{\Psi} = \left( \sum_{i=1}^m u_i B_i \right) \Psi$$

**where the  $B_i$  are Hermitian control operators.**

**Assume**

- $H_1$  is an invariant subspace for  $B_1$
- the system is unit vector controllable on  $H_1$  using only  $B_1$

# Finite controllability theorem (cont'd)

**If**

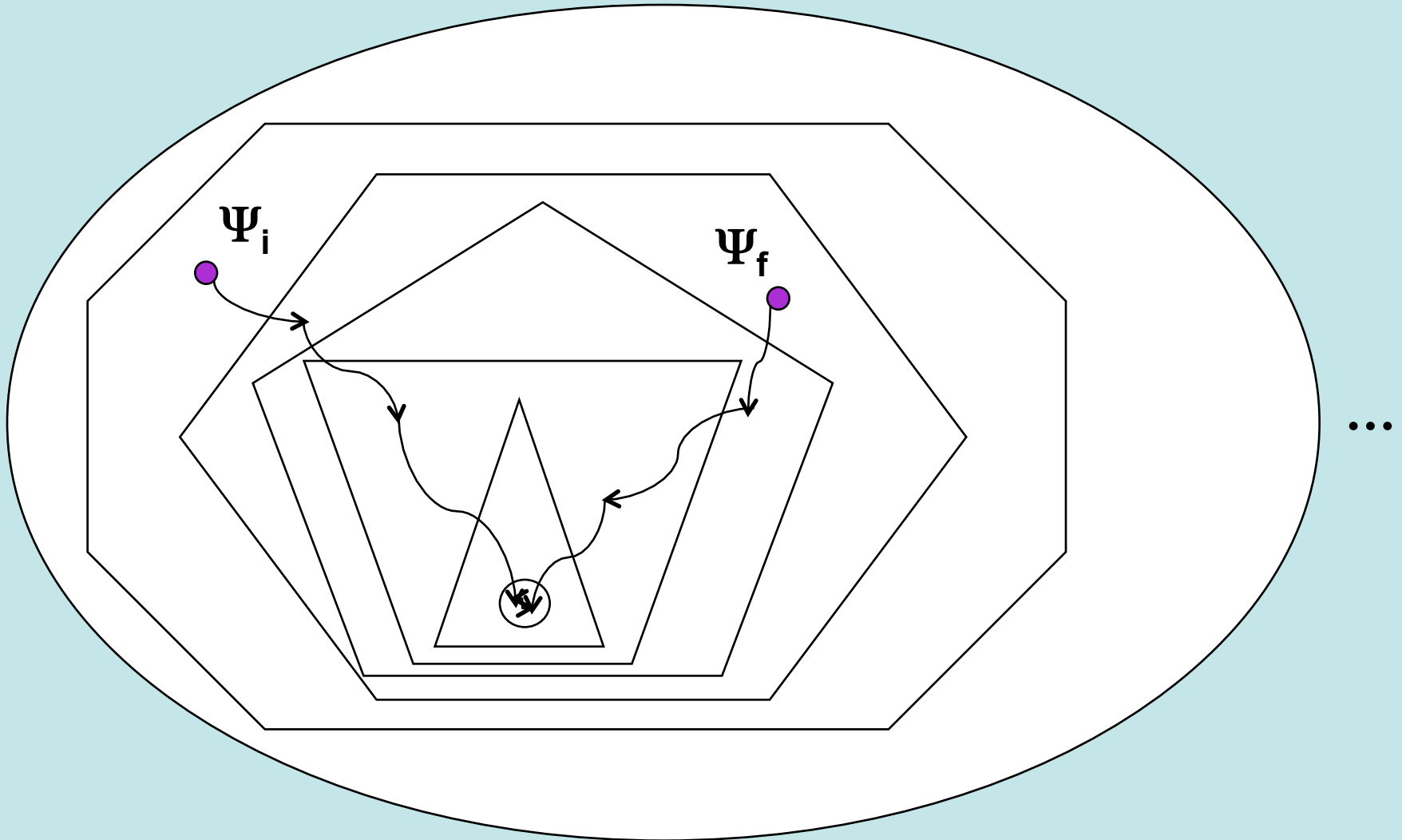
- for each  $H_\alpha; \alpha \neq 1$  there is a  $B_\alpha$  that leaves  $H_\alpha$  invariant, and
- for any unit vector in  $H_\alpha$  the orbit generated by  $\exp(iB_\alpha)$  contains a point in one of the lower dimensional subspaces  $H_\beta$

**then any unit vector in any of the  $H_i$  can be steered to any other unit vector in any other  $H_j$  using a finite number of piecewise constant controls.**

**Bloch, Brockett, Rangan, [quant-ph/0608075](https://arxiv.org/abs/quant-ph/0608075); ITAC 55, 1797 (2010)**

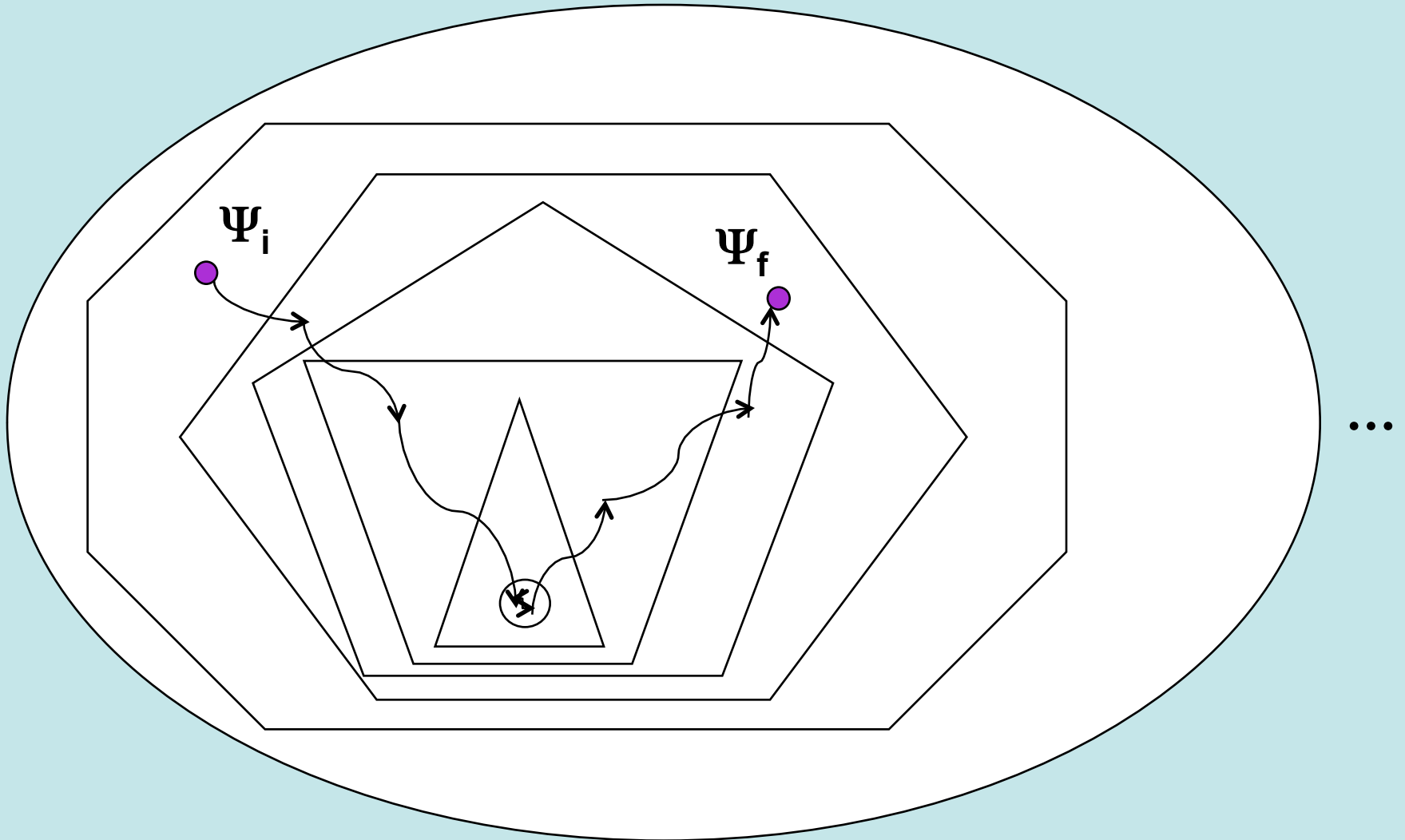
# Finite controllability

H



# Explicit scheme

H



# E.g.: trapped-ion quantum states

Physical system:

Trapped-ion

Model:

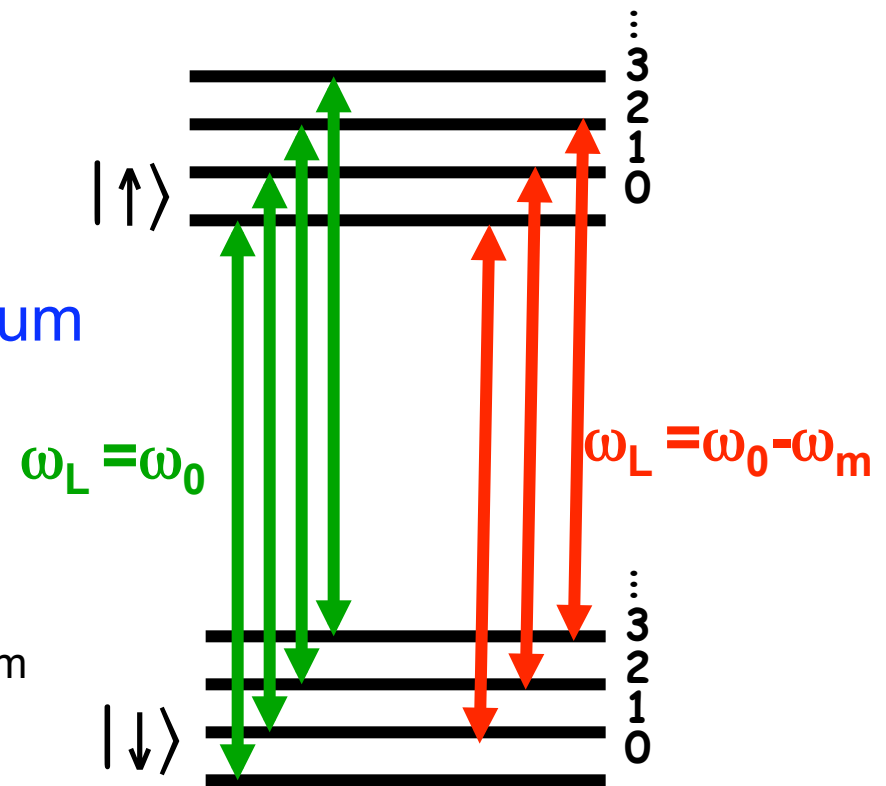
Spin  $\frac{1}{2}$  system coupled to quantum harmonic oscillator

Field-free Hamiltonian:

$$H_0 = (\frac{1}{2})\omega_0\sigma_z + \omega_m a^\dagger a \quad ; \quad \omega_0 \gg \omega_m$$

Control field:

$$E(\xi, t) = \hat{e} E(t)\cos(k\xi - \omega_L t)$$



Field-free eigenstates transitively connected by two resonant, monochromatic fields



# Mathematical formulation

## Interaction Hamiltonian:

$$\begin{aligned} H_I &= (1/2) \mu \sigma_x E(t) \cos(k\xi_0(a+a^\dagger) - \omega_L t) \\ &= \sigma_x \Omega(t) \cos(\eta(a+a^\dagger) - \omega_L t) \end{aligned}$$

Lamb-Dicke parameter  $\eta = k\xi_0$

## Go into rotating frame, make RWA

Non-zero matrix elements in control Hamiltonian:

$$\langle S, n | H_I | S', m \rangle$$

**Carrier** ( $\omega_L = \omega_0$ ):

$$|\downarrow n\rangle \leftrightarrow |\uparrow n\rangle \sim L_n(\eta^2)$$

**First red sideband** ( $\omega_L = \omega_0 - \omega_m$ ):

$$|\downarrow n\rangle \leftrightarrow |\uparrow n-1\rangle \sim i L_n^{(1)}(\eta^2)$$

D. J. Wineland *et al.*, *J. Res. Natl. Inst. Stand. Technol.* 103, 259 (1998).

D. Leibfried, R. Blatt, C. Monroe, D. Wineland. *Rev. Mod. Phys.* 75, 281 (2003).

# Infinite Lie algebra

$$i\dot{\Psi} = (uB_c + vB_r)\Psi = (H_c + H_r)\Psi$$

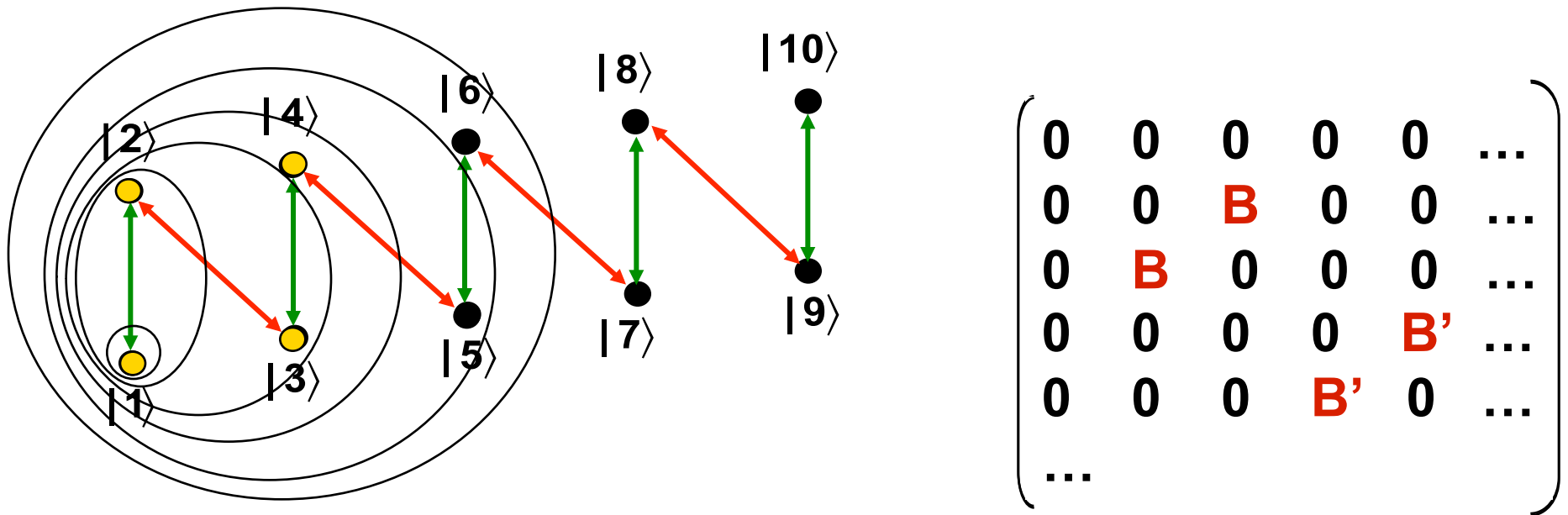
$$\begin{pmatrix} 0 & \mathbf{A} & 0 & 0 & 0 & \dots \\ \mathbf{A} & 0 & 0 & 0 & 0 & \dots \\ 0 & 0 & 0 & \mathbf{A}' & 0 & \dots \\ 0 & 0 & \mathbf{A}' & 0 & 0 & \dots \\ 0 & 0 & 0 & 0 & 0 & \dots \\ \dots & & & & & \end{pmatrix}$$

$$\begin{pmatrix} 0 & 0 & 0 & 0 & 0 & \dots \\ 0 & 0 & \mathbf{B} & 0 & 0 & \dots \\ 0 & \mathbf{B} & 0 & 0 & 0 & \dots \\ 0 & 0 & 0 & 0 & \mathbf{B}' & \dots \\ 0 & 0 & 0 & \mathbf{B}' & 0 & \dots \\ \dots & & & & & \end{pmatrix}$$

$$\begin{aligned} & \exp(-iH\Delta t) \\ &= \exp(-i(H_c + H_r)\Delta t) \\ &= \exp(-iH_c\Delta t) \cdot \exp(-iH_r\Delta t) \\ &\quad \cdot \exp(-1/2[H_c, H_r](\Delta t)^2) \\ &\quad \cdot \exp(1/12[H_c, [H_c, H_r]](\Delta t)^3) \\ &\quad \cdot \exp(1/12[[H_c, H_r], H_r](\Delta t)^3) \dots \end{aligned}$$

**Lie algebra is  $\infty$ -D**  
 The alternate application of control fields removes a chirp instability in unitary flows. (Brockett, Rangan, & Bloch, CDC 2003)

# Finite controllability of trapped-ion



Reachable set includes superpositions of finite numbers of eigenstates.

Approximate controllability of trapped-ion quantum states  
(BBR, quant-ph/0608075, ITAC 2010)

Also see: Ervedoza-Puel, Ann. I. H. Poincaré 26, 2111 (2009)

# Lamb-Dicke limit

Transition matrix elements:  $\langle \Phi_1 | H_I | \Phi_2 \rangle$

Carrier:

$$|\downarrow n\rangle \leftrightarrow |\uparrow n\rangle \sim L_n(\eta^2)$$

First red sideband:

$$|\downarrow n\rangle \leftrightarrow |\uparrow n-1\rangle \sim i L_n^{(1)}(\eta^2)$$

In practical implementations, finding the control pulses is difficult (annoying) since the matrix elements depend on Laguerre polynomials

Lamb-Dicke limit (LDL): (motional cooling)

$$\xi_0 \ll \lambda, \eta \ll 1$$

Carrier:

$$|\downarrow n\rangle \leftrightarrow |\uparrow n\rangle \sim 1$$

First red sideband:

$$|\downarrow n\rangle \leftrightarrow |\uparrow n-1\rangle \sim \sqrt{n}$$

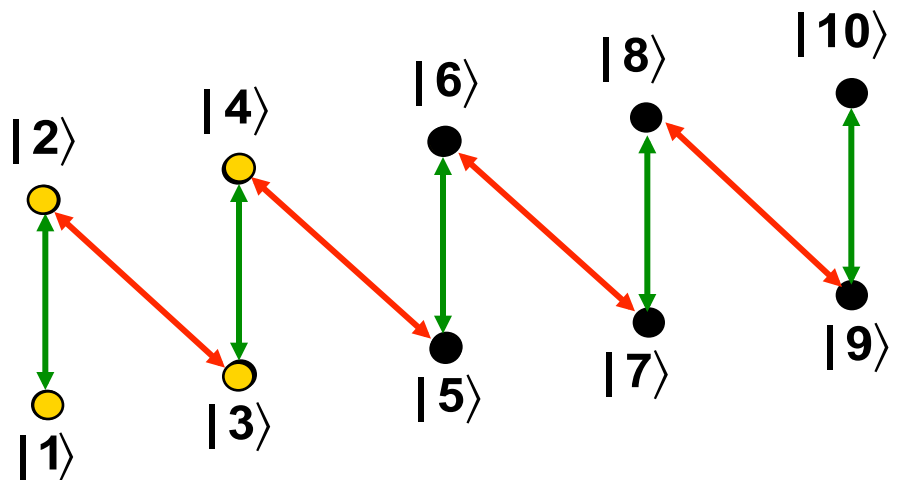
# Explicit control schemes

In LDL: Law-Eberly scheme, PRL 76, 1055 (1996)  
Kneer-Law scheme, PRA 57, 2096 (1998)

Beyond LDL: Wei, Liu & Nori, PRA 70, 63801 (2004)

Aim: Start from ground state and create a finite superposition of trapped-ion energy eigenstates

Method: reverse engineer



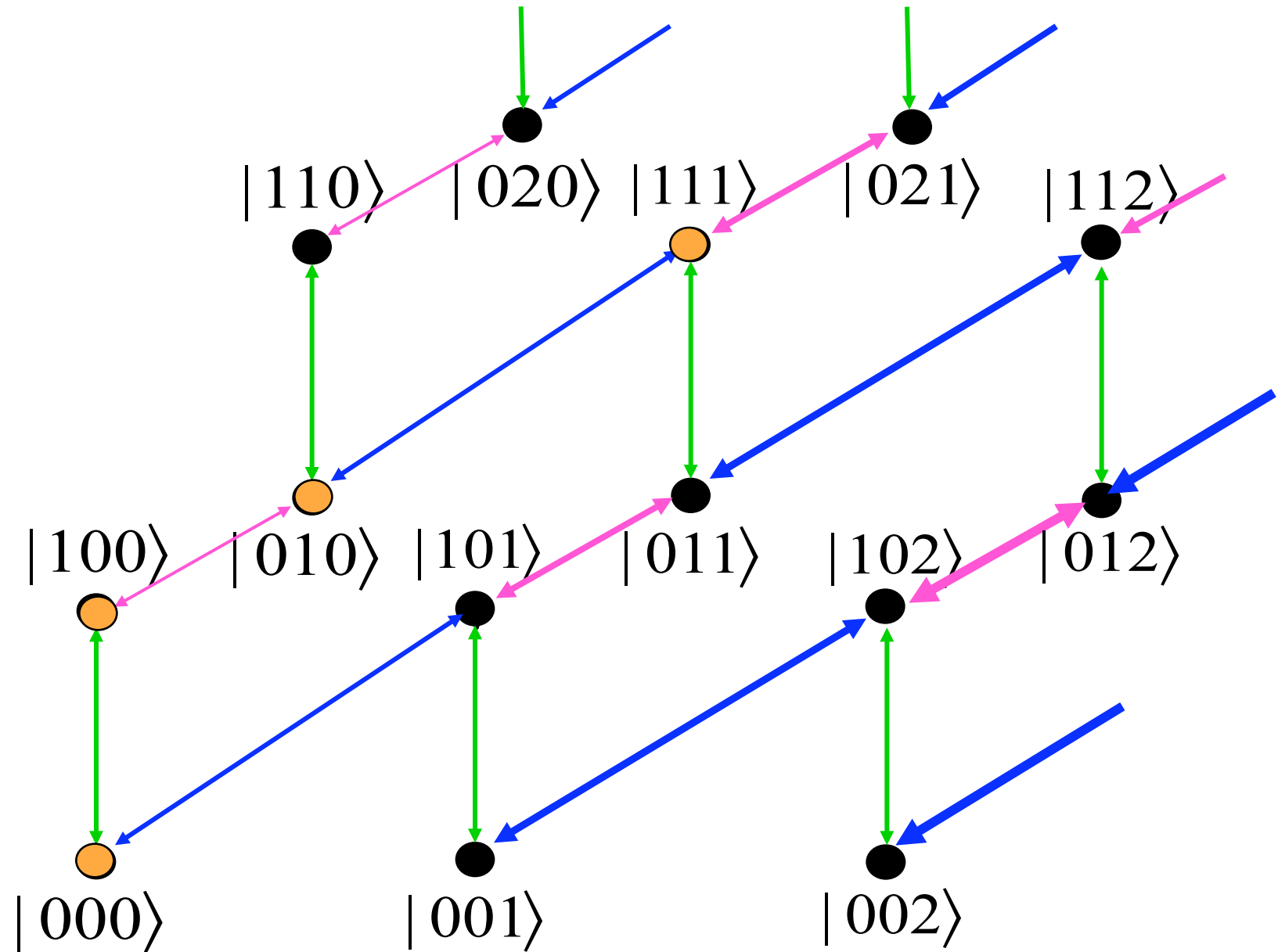
Experimental demonstration:

(LDL) Ben-Kish et al., Phys. Rev. Lett. 90, 037902 (2003).

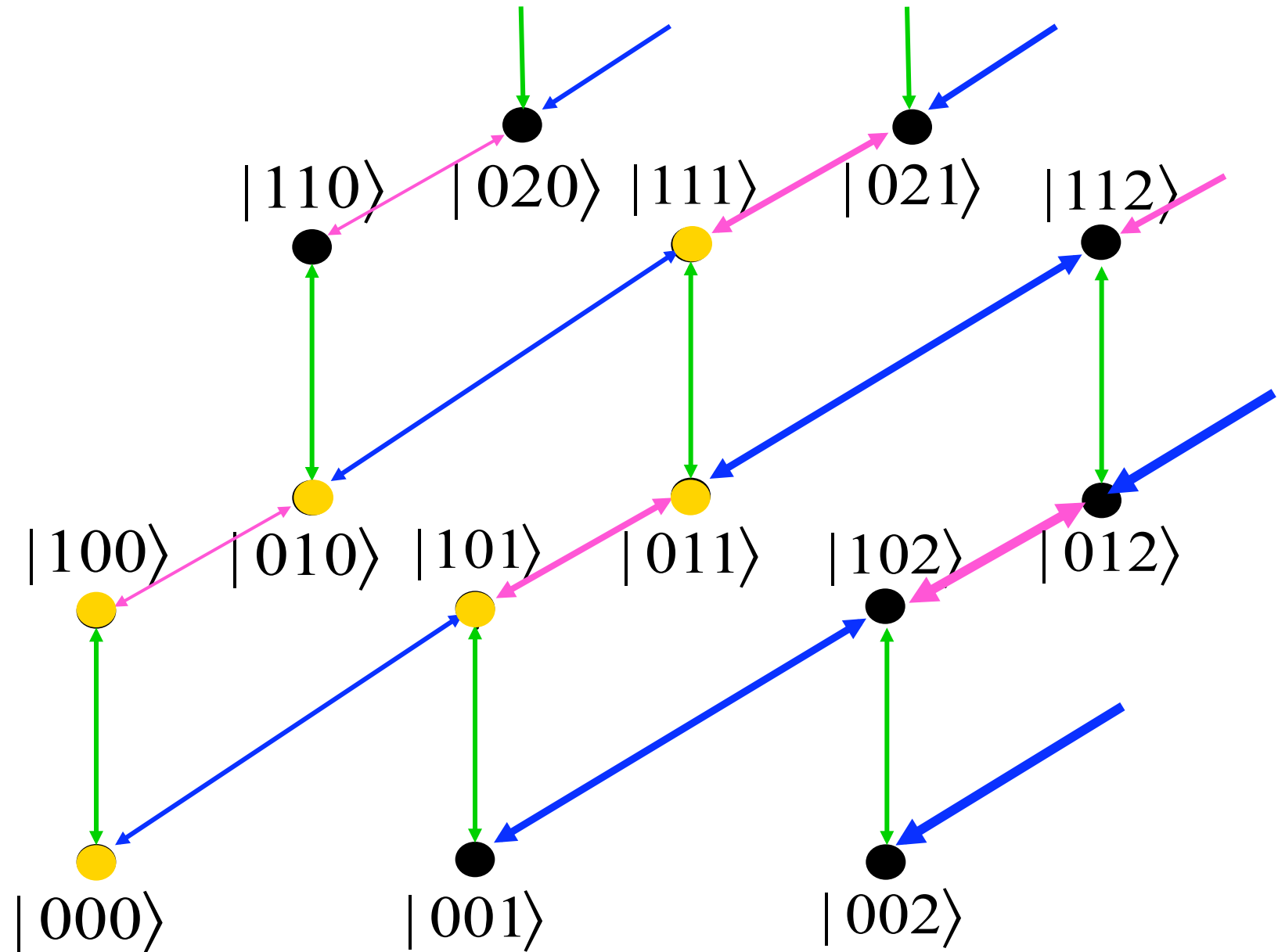
Without LDA (demonstration of principle):

B.E. King et al., Phys. Rev. Lett. 81, 3631 (1998).

# E.g.: spin-1/2 with 2 QHO's



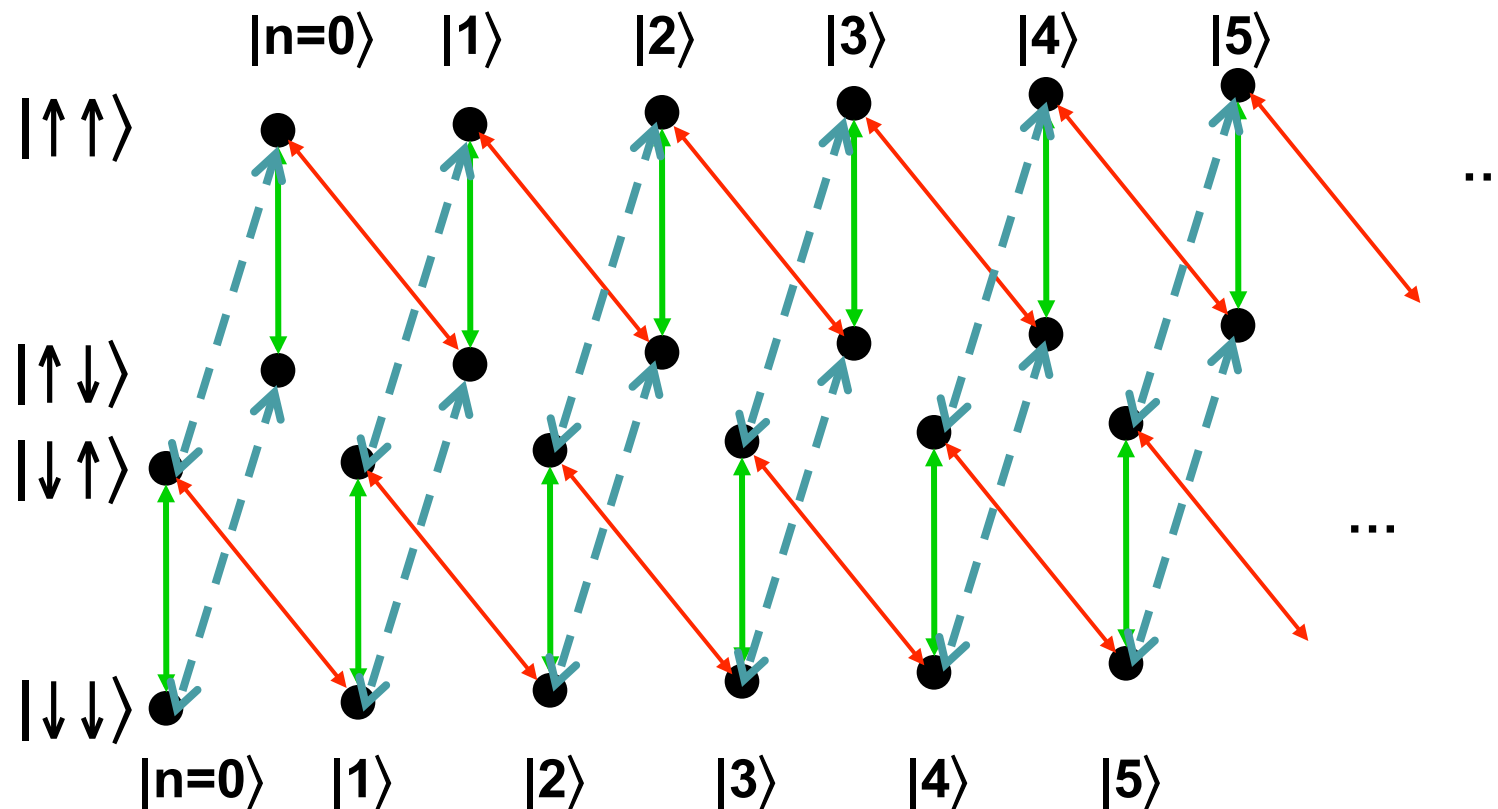
# BUT - No finite controllability



# E.g.: Two trapped-ions

Two spin  $\frac{1}{2}$  systems coupled to a quantum harmonic oscillator

Assume we have independent controls  $u^{(1)}$ ,  $u^{(2)}$  and  $v^{(1)}$



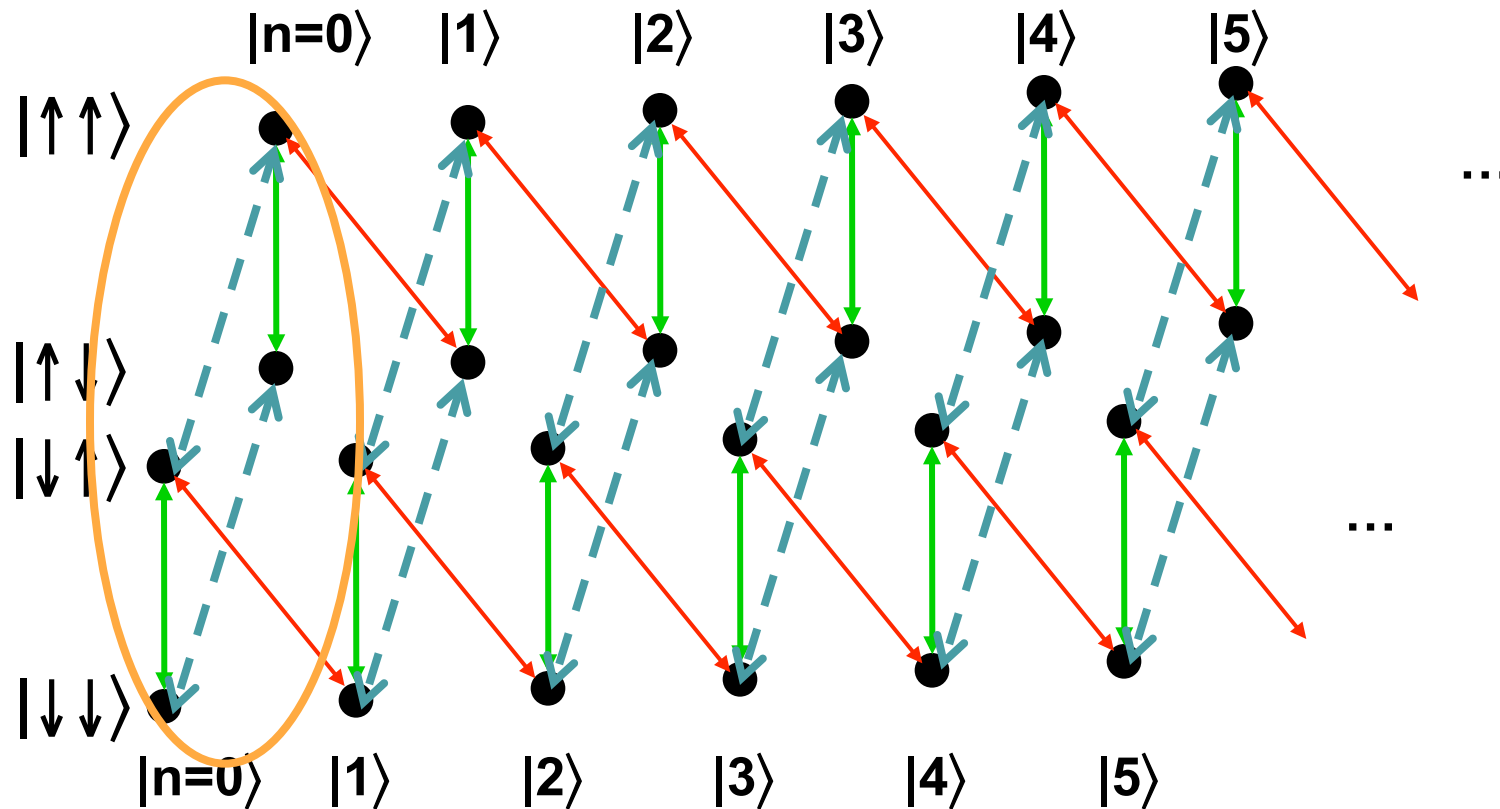
Have eigenstate controllability, can also make specific types of states: e.g.: Turquette et al., Phys. Rev. Lett. 81, 3631 (1998).



# E.g.: Two trapped-ions

Two spin  $\frac{1}{2}$  systems coupled to a quantum harmonic oscillator

Assume we have independent controls  $u^{(1)}$ ,  $u^{(2)}$  and  $v^{(1)}$

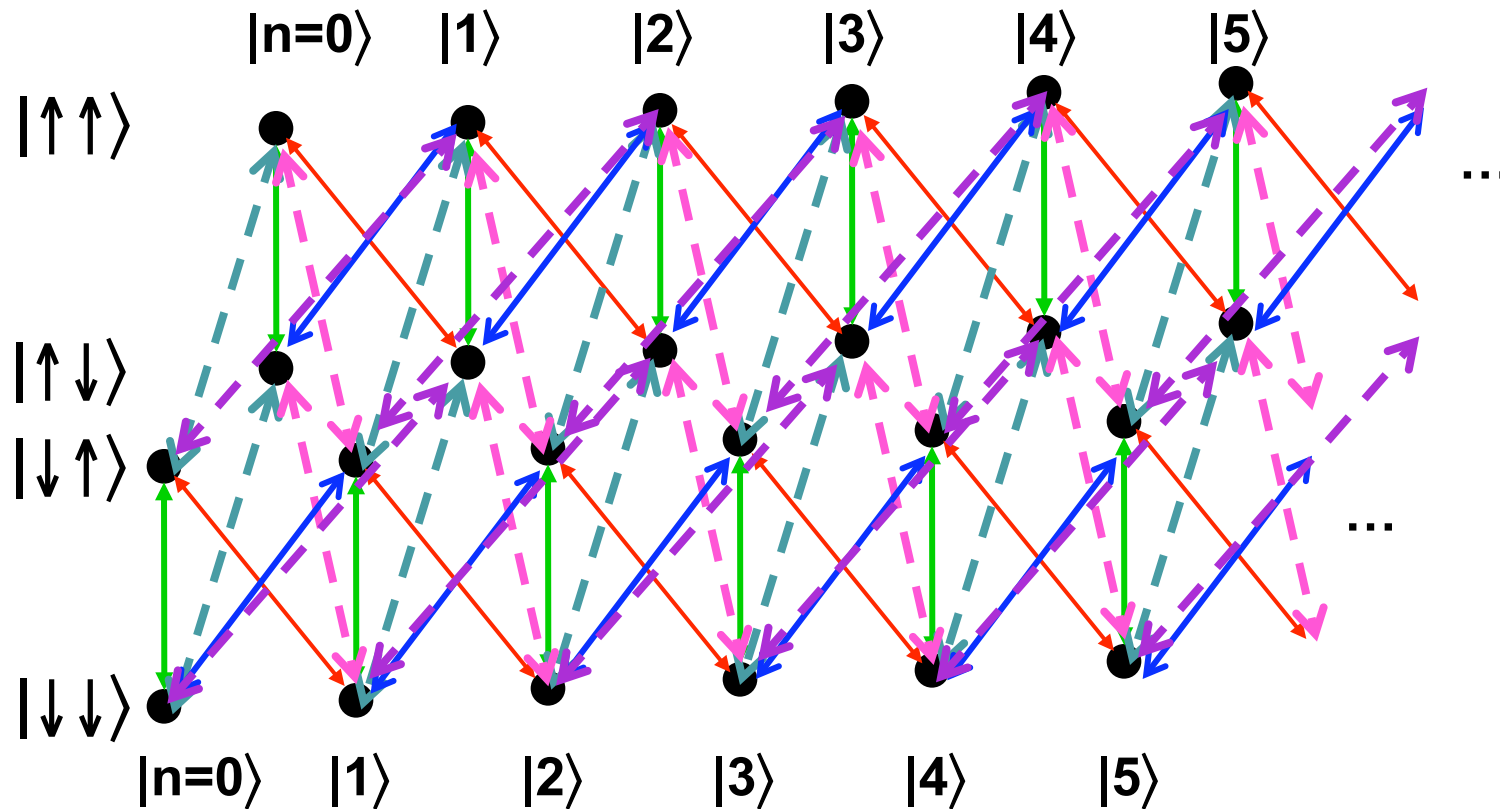


...but does not satisfy conditions for finite controllability.

# Approximate controllability?

Two spin  $\frac{1}{2}$  systems coupled to a quantum harmonic oscillator

Assume we have 6 independent controls, still not fin. cont.



How to start from  $\frac{1}{2}|\downarrow\uparrow 0\rangle + \frac{\sqrt{3}}{2}|\uparrow\uparrow 0\rangle + \frac{\sqrt{3}}{2}|\uparrow\downarrow 0\rangle$  and get  $|\downarrow\downarrow 0\rangle$  using a single non-zero control at a time?

# Summary

## **Finite controllability**

-Approx. controllability of spin-half particle coupled to an harmonic oscillator

## **No finite controllability for**

- Spin-half particle coupled to two harmonic oscillators
- Two spin-half particles coupled to an harmonic oscillator
  - Conjecture: the latter system is not approximately controllable using sequential fields

# Acknowledgements

Students: D. Rooney, S. Mirzaee, A. Torabi, D. Travo, M. Sheikh, J. Donohue, M. Tywoniuk, pdf T. Cheng

Collaborators: Silvia Mittler @ UWO, Tony Bloch @ U.Mich, Roger Brockett @Harvard

**BiopSys: NSERC  
Strategic Network on  
Bioplasmonic Systems**



Ontario Centres of  
Excellence



Windsor Regional Cancer Centre  
Centre Régional De Cancérologie De Windsor  
patient centred excellence



UNIVERSITY OF  
WINDSOR

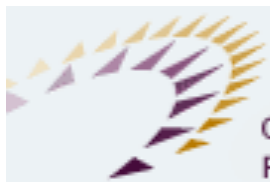


Ontario

OAML  
OAML



MITACS



Canada Foundation for Innovation  
Fondation canadienne pour l'innovation



NSERC  
CRSNG



ICIP  
CIPI