A Quasi-Stationary Distribution Approach to Transient Dynamics

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Retreat for Young Researchers in Probability and areas of Application $${\rm BIRS}$$

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Collaborators

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Stochastic chemical reaction models

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QSDs and transient states

Conclusion and vision

Illustration of transient dynamics (from R. Capeáns, J. Sabuco and M. A. F. Sanjuán, DCDS-B, 2018)



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Illustration of the flow at a circular pier (from B.W. Melville and S.E. Coleman, Bridge Scour Water, 2000)



Three-species food chain model with cooperative hunting (from J. Duarte, C. Januário, N. Martins and J. Sardanyés, Chaos, 2009):

$$(\text{Resource}) \quad \dot{R} = R\left(1 - \frac{R}{K}\right) - \frac{x_c y_c CR}{R + R_0},$$

$$(\text{Consumer}) \quad \dot{C} = -x_c C + \frac{x_c y_c CR}{R + R_0} - \phi(P) \frac{x_p y_p C}{C + C_0},$$

$$(\text{Predator}) \quad \dot{P} = -x_p P + \phi(P) \frac{x_p y_p C}{C + C_0},$$

where

$$\phi(P) = (1 - \sigma)P + \sigma P^2.$$

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$$\begin{aligned} &(\text{Resource}) \quad \dot{R} = R\left(1 - \frac{R}{K}\right) - \frac{x_c y_c CR}{R + R_0}, \\ &(\text{Consumer}) \quad \dot{C} = -x_c C + \frac{x_c y_c CR}{R + R_0} - \phi(P) \frac{x_p y_p C}{C + C_0}, \\ &(\text{Predator}) \quad \dot{P} = -x_p P + \phi(P) \frac{x_p y_p C}{C + C_0}, \end{aligned}$$

where

$$\phi(P) = (1 - \sigma)P + \sigma P^2.$$

 σ ∈ [0, 1] is a measure of the degree of cooperation inside the population of predators.



Figure: (a) $\sigma = 0$, McCann-Yodzis model; (b) (c) $0 < \sigma \ll 1$

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More examples:

Dispersion of volcanic ash;

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- Dispersion of volcanic ash;
- Pollutant spreading in the atmosphere;

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Process of decision making.

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Traditional approaches:

normally hyperbolic invariant manifold;



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identifying the time scales;

More examples:

- Dispersion of volcanic ash;
- Pollutant spreading in the atmosphere;
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Traditional approaches:

normally hyperbolic invariant manifold;



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- identifying the time scales;
- perturbing chaotic systems + measure of complexity (Lyapunov exponent, entropy, dimension, etc.).

Chemical reactions:

$$A + X \xrightarrow[k_{-1}]{k_{-1}} 2X, \quad X \xrightarrow{k_2} C,$$

• reaction rates: k_1 , k_{-1} , k_2 ,

• open system, the number of A molecules held fixed n_A ,

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scaling parameter V = volume × Avogadro number.

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$$A + X \xrightarrow[k_{-1}]{k_{-1}} 2X, \quad X \xrightarrow{k_2} C,$$

• reaction rates: k_1 , k_{-1} , k_2 ,

- open system, the number of A molecules held fixed n_A ,
- scaling parameter V = volume × Avogadro number.

Continuous-time Markov jump process:

• X_t^V : the process counting the number of X molecules,

• $\frac{X_t^{V}}{V}$: the process for the concentration.

Stochastic chemical reaction models Master equation for $\frac{X_t^V}{V}$:

 $\frac{d}{dt}p(t,\frac{n}{V})=\lambda_{n-1,V}p(t,\frac{n-1}{V})-(\lambda_{n,V}+\mu_{n,V})p(t,\frac{n}{V})+\mu_{n+1,V}p(t,\frac{n+1}{V}),$

where $p(t, \frac{n}{V}) = \mathbb{P}\left[\frac{X_t^V}{V} = \frac{n}{V}\right]$, and $\lambda_{n,V}, \mu_{n,V}$ are determined by law of mass action.



$$\frac{k_{-1}n(n+1)}{V} = \frac{k_{-1}}{V^{2-1}} 2! \binom{n+1}{2},$$

$$A + X \xrightarrow[k_{-1}]{k_{-1}} 2X, \quad X \xrightarrow{k_2} C.$$

Mean-field approximation of $\frac{X_t^V}{V}$:

$$\dot{x} = b(x), \quad x \in [0,\infty)$$

where

$$b(x) = k_1 x_A x - k_{-1} x^2 - k_2 x = k_{-1} x \left(\frac{k_1 x_A - k_2}{k_{-1}} - x \right).$$

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x_A is concentration of *A* molecules (^{n_A}/_V → *x_A* as *V* → ∞).
 Assume *k*₁*x_A* > *k*₂. Positive equilibrium *x_e* := ^{k₁x_A-k₂}/_{k-1}.

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- Assume $k_1 x_A > k_2$. Positive equilibrium $x_e := \frac{k_1 x_A k_2}{k_{-1}}$.

Dynamics:

$$x(t)
ightarrow x_e$$
 exponentially as $t
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Diffusion approximation of $\frac{X_t^V}{V}$:

$$dx = b(x)dt + \epsilon \sqrt{a(x)}dW_t, \quad x \in [0,\infty),$$

where $\epsilon = \frac{1}{\sqrt{V}}$ and

$$a(x) = k_1 x_A x + k_{-1} x^2 + k_2 x.$$

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- δ_0 is the unique stationary distribution.

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 - Extinction time T_{ϵ} satisfies $\mathbb{E}[T_{\epsilon}] < \infty$.
- Stochastic stability: small noise stabilizes/de-stabilizes the unstable/stable equilibrium.

Two-scale dynamics of the diffusion approximation

$$dx = b(x)dt + \epsilon \sqrt{a(x)}dW_t,$$

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Eventual extinction dynamics.

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- Eventual extinction dynamics.
- Transient dynamics: stay near x_e for a (long) finite time period.

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$$dx = b(x)dt + \epsilon \sqrt{a(x)}dW_t,$$

- Eventual extinction dynamics.
- Transient dynamics: stay near x_e for a (long) finite time period.
 - Large deviation. For each T > 0,

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\mathsf{SDE}_{\epsilon} \approx \mathsf{ODE} \text{ on } [0, T]
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for all $0 < \epsilon \ll 1$.

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$$dx = b(x)dt + \epsilon \sqrt{a(x)}dW_t,$$

- Eventual extinction dynamics.
- Transient dynamics: stay near x_e for a (long) finite time period.
 - Large deviation. For each T > 0,

$$SDE_{\epsilon} \approx ODE \text{ on } [0, T]$$

for all $0 < \epsilon \ll 1$.

▶ Beyond large deviation. For each 0 < $\epsilon \ll 1$, there is T_{ϵ} such that

 $SDE_{\epsilon} \approx ODE \text{ on } [0, T_{\epsilon}],$ $SDE_{\epsilon} \not\approx ODE \text{ on } (T_{\epsilon}, \infty).$

Questions:

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- How long do solutions stay with the transient state?

Quasi-stationary distribution: a QSD μ_{ϵ} is a probability measure on $(0, \infty)$ s.t.

$$\mathbb{P}_{\mu_{\epsilon}}\left[x_{\epsilon}(t)\inullet|\mathcal{T}_{\epsilon}>t
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principal eigen-pair:

$$\mathcal{L}^{\epsilon} u_{\epsilon} = -\lambda_{\epsilon} u_{\epsilon},$$

where
$$\lambda_\epsilon > 0$$
 and $\mathcal{L}^\epsilon u = rac{\epsilon^2}{2} (au)_{xx} - (bu)_x;$

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Problem: λ_{ϵ} Vs η_{ϵ} .

- $\lambda_{\epsilon} > \eta_{\epsilon} \implies$ extinction ahead of QSD;
- $\lambda_{\epsilon} < \eta_{\epsilon} \implies$ QSD ahead of extinction.

Theorem (Z. Shen, S. Wang and Y. Yi, 2019)

• For each $\mathcal{O} \subset \subset (0,\infty) \setminus \{x_e\}$, there are $\gamma_{\mathcal{O}}$ and $\epsilon_{\mathcal{O}}$ s.t.

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Commutative diagram:



Conclusion:

- use QSDs to capture the transient states;
- study the concentration of QSDs;
- obtain results that are compatible with the intuition.

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computation: could be more stable and accurate.

Thank you for your attention!

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