Gap Probabilities in tiling models and discrete Painlevé equations

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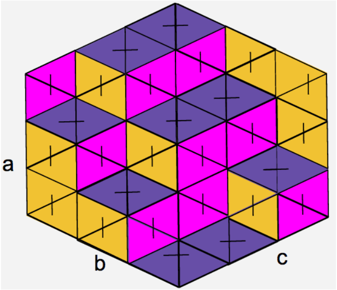
8. As an example, we consider the computation of *gap probabilities* in a generalized tiling problem (Alisa Knizel's work).

Probabilistic Model: q-Distributions on Boxed Plane Partitions

- Models of a random surfaces: *boxed plane partition* (*lozenge tiling* of a hexagon).
- Consider tilings of an $a \times b \times c$ hexagon ($a, b, c \geq 1$) by three types of lozenge tilings (obtained by gluing together two adjacent triangles of a regular triangular grid).



Denote the set of all possible such tilings by $\Omega_{a\times b\times c}$. Equip this set with a probability measure, where, for $\mathcal{T} \in \Omega_{a\times b\times c}$,

$$P(\mathcal{T}) = \frac{w(\mathcal{T})}{Z(a,b,c)}, \text{ where } w(\mathcal{T}) = \prod_{\diamond \in \mathcal{T}} w(\diamond),$$

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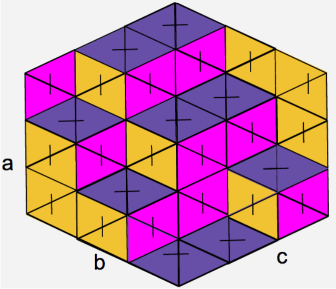

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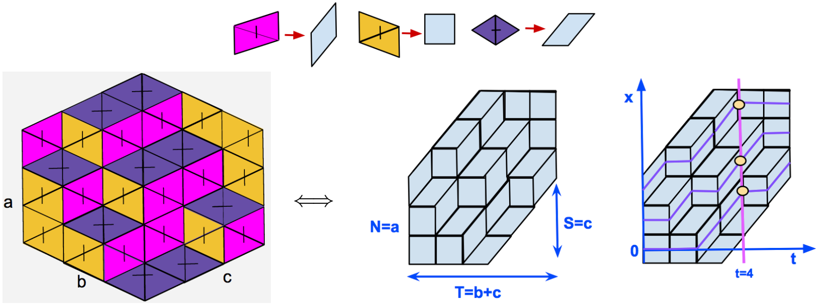
- Originally, the most studied distribution was uniform, $w(\diamond) = 1$. In 2009, A. Borodin, V. Gorin, and E. Rains introduced a far-reaching generalization of this model with a very general *elliptic weight* and (complex) parameters u_1, u_2, p, q:

$$w(\diamond) = w(\diamond_{i,j}) = \frac{(u_1u_2)^{1/2}q^{j-1/2}\theta_p(q^{2j-1}u_1u_2)}{\theta_p(q^{j-3i/2-1}, q^{j-3i/2}u_1, q^{j+3i/2-1}, q^{j+3i/2}u_2)},$$

where $\theta_p(x) = \prod_{i=0}^{\infty}(1-p^ix)(1-p^{i+1}/x)$ and $\theta_p(a,b,c,\dots) = \theta_p(a)\theta_p(b)\theta_p(c)\dots$.

Change variables to $N = a$, $T = b + c$, $S = c$ and interpreting plane partitions as nonintersecting paths via an affine transformation



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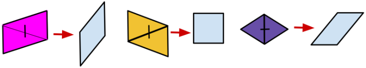
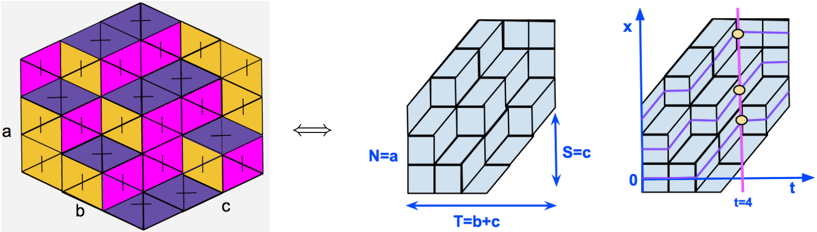
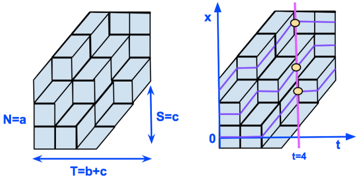




- The most general elliptic weight $w(\diamond) = \dfrac{(u_1 u_2)^{1/2} q^{j-1/2} \theta_p(q^{2j-1} u_1 u_2)}{\theta_p(q^{j-3i/2-1}, q^{j-3i/2} u_1, q^{j+3i/2-1}, q^{j+3i/2} u_2)}$ corresponds to certain biorthogonal functions (not polynomials).
- The most general orthogonal polynomial case is the limit $p \to 0$, $u_1 = O(\sqrt{p})$, $u_2 = O(\sqrt{p})$, $u_1 u_2 = p\kappa^2 q^{-S}$ with the *q-Racah weight* $w(\diamond) = \kappa q^{j-(S+1)/2} - \dfrac{1}{\kappa q^{j-(S+1)/2}}$.
- Taking the limit with $\kappa \to 0$ (with appropriate rescaling) gives q-Hahn weights $w(\diamond) = q^{-j}$.

View the boxed partition model as the non-intersecting paths model; equip it with the q-Hahn weight $w(\diamond) = q^{-j}$. Fix a section t. Let the coordinates of the nodes be $C(t) = (x_1, \ldots, x_N)$.



Gap Probabilities (the q-Hahn case)

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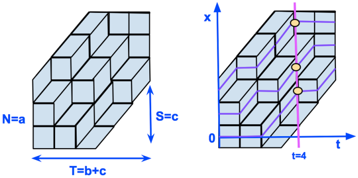

Theorem (Borodin, Gorin, Rains (2009))

$$\mathrm{Prob}\{C(t) = (x_1, \ldots, x_N)\} = const \cdot \prod_{0 \le i < j \le M} (q^{-x_i} - q^{-x_j})^2 \prod_{i=1}^{N} w(x_i),$$

where $w(x)$ is the weight function of the q-Hahn polynomial ensemble up to a factor not depending on x.

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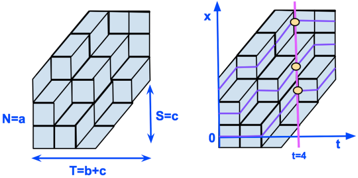

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Gap probability

The one-interval gap probability function D_s^N is

$$D_s^N = \mathrm{Prob}[\max\{x_i\} < s].$$

Theorem (Knizel (2016), q-Volume case)

The gap probability D_s^N for the q-Hahn ensemble can be computed recursively

$$D_s^N = \frac{(D_{s-2}^N)^2}{D_{s-1}^N} \frac{(r_{s-1}w - qvz_1z_2)(r_s w - quz_1z_2)(t_{s-1} - qz_1)(t_{s-1} - qz_2)}{uvz_1z_2(qz_1 - z_3)(qz_1 - z_5)(qz_2 - z_4)(qz_2 - z_6)},$$

where the sequence (r_s, t_s) satisfies the recursion (equivalent to after some change of parameters)

$$(r_s t_{s-1} + 1)(r_{s-1} t_{s-1} + 1) = \frac{z_1z_2(t_{s-1} - z_3)(t_{s-1} - z_4)(t_{s-1} - z_5)(t_{s-1} - z_6)}{z_3z_4z_5z_6(qt_{s-1} - z_1)(qt_{s-1} - z_2)},$$

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The parameters $u, v, w, z_1, \ldots, z_6$ and the initial conditions are explicitly computed in terms of α, β, q, s. The above recursion coincides with the q-P $\left(A_2^{(1)}/E_6^{(1)}\right)$ of (**KNY**) *after some change of parameters.*

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- This relation is obtained through the DRHP approach, that can be interpreted as describing isomonodromy deformations of a q-connection.
- Moduli space of such connections turn out to coincide with Sakai's q-Painlevé surfaces.
- Thus, the isomonodromy deformations of connections are maps in this q-Painlevé family, and hence should be given by q-P equations.

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How to identify them?

DRHP for q-Hahn

Theorem (Borodin-Boyarchenko (2002))

Fix $\operatorname{card}(\mathfrak{X}) > k > 0$ *and set* $w(\psi) = \begin{bmatrix} 0 & w(\psi) \\ 0 & 0 \end{bmatrix}$. *For any* $s \geq k$ *there exists unique analytic function* $m_s(\psi) : \mathbb{C} \setminus \mathfrak{N}_s \to \operatorname{Mat}(\mathbb{C}, 2)$ *with simple poles at points in* $\mathfrak{N}_s = \{x_0, \dots, x_{s-1}\}$ *such that*

$$\operatorname{Res}_{\psi = x} m_s(\psi) = \lim_{\psi \to x} m_s(\psi) w(\psi), \quad x \in \mathfrak{N}_s;$$

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Introduce matrix

$$A_s(z) = m_s(q^{-1}z) A_0(z) m_s^{-1}(z), \qquad \text{where} \quad A_0(z) = \begin{bmatrix} \dfrac{q w(x+1)}{w(x)} & 0 \\ 0 & 1 \end{bmatrix} \text{ with } z = q^{-x}.$$

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In the q-Hahn case,

$$\frac{q w(x+1)}{w(x)} = \frac{(z - \alpha q) \cdot (z - q^{-M})}{\alpha\beta (z - q) \cdot (z - \beta^{-1} q^{-M})}.$$

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- Evolution $(t_s, p_s) \to (t_{s+1}, p_{s+1})$ is described by (**after some adjustments**) q-P$(A_2^{(1)})$.

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- For any s matrix element a_{21}^s has a unique zero. Denote it by t_s and $a_{11}^s(t_s)$ by p_s.
- Evolution $(t_s, p_s) \to (t_{s+1}, p_{s+1})$ is described by (**after some adjustments**) q-P$(A_2^{(1)})$.

Structure of a generic $A_s(z)$ of type $\lambda = (z_1, \ldots, z_6; u, v, w, w; 3)$

$$A(z) = \begin{bmatrix} a_{11}(z) & a_{12}(z) \\ a_{21}(z) & a_{22}(z) \end{bmatrix}, \quad A(0) = \begin{bmatrix} w & 0 \\ 0 & w \end{bmatrix},$$

where $\deg(a_{11}) \leq 3$, $\deg(a_{12}) \leq 2$, $\deg(a_{21}) \leq 2$, $\deg(a_{22}) \leq 3$ and

$$\det A(z) = uv(z - z_1)(z - z_2)(z - z_3)(z - z_4)(z - z_5)(z - z_6)$$

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Parameter evolution

When $A_s(z) \to A_{s+1}(z)$, the parameters evolve as

$$(z_1^s, z_2^s, \ldots, z_6^s, u_s, v_s, w_s) \to (z_1^{s+1}, z_2^{s+1}, \ldots, z_6^{s+1}, u_{s+1}, v_{s+1}, w_{s+1})$$

with $z_2^{s+1} = qz_2^s$, $z_4^{s+1} = qz_4^s$, $w_{s+1} = qw_s$, and $z_i^{s+1} = z_i^s$ for $i \neq 2, 4$.

Weight degenerations and Sakais Classification scheme for Discrete Painlevé equations

The main goal of this project is to both find a way to extend the results from the q-Hahn case to a more general q-Racah case, and also to see how it fits the degeneration cascade in Sakai's classification scheme for discrete Painlevé equations.

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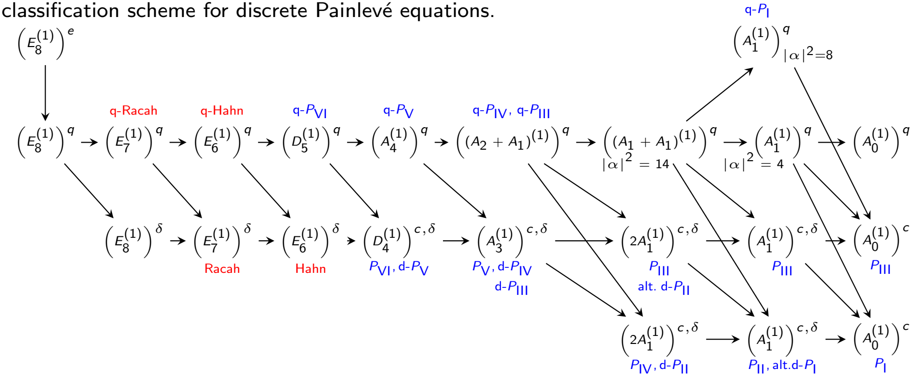

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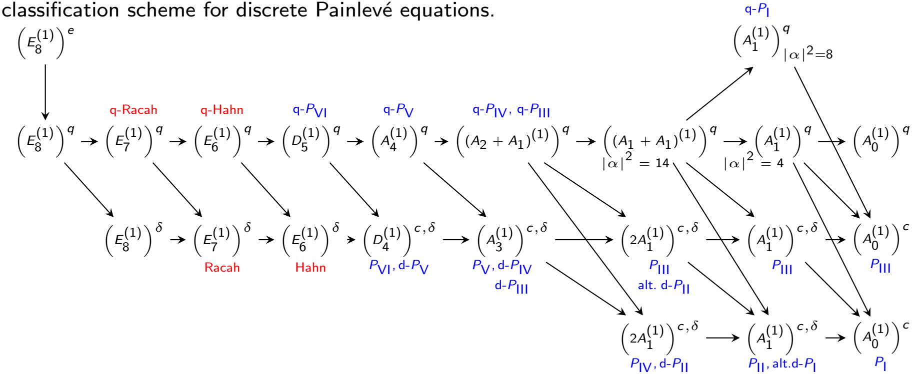
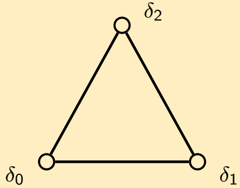


Every discrete Painlevé equation is a discrete dynamical system given by a non-homogeneous birational automorphism of $\mathbb{P}^1 \times \mathbb{P}^1$. It is resolved by blowing up $\mathbb{P}^1 \times \mathbb{P}^1$ at eight points, and becomes a flow on a family of such surfaces. Configuration of blowup points is encoded by an affine Dynkin diagram. Its "dual" affine Dynkin diagram encodes the affine Weyl symmetry group of the family (above) and Discrete Painlevé equation is equivalent to a *translation* in its lattice.

$A_2^{(1)}$ surface model



$$\delta_0 = H_f + H_g - F_1 - F_2 - F_3 - F_4$$
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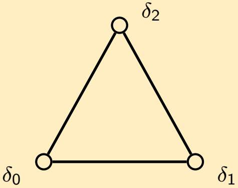

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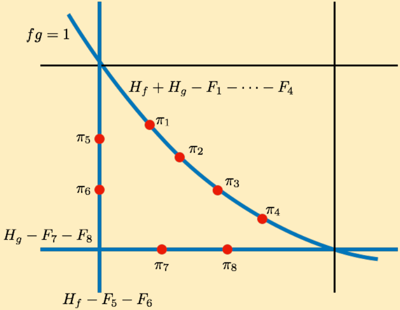

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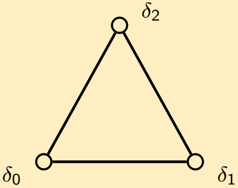

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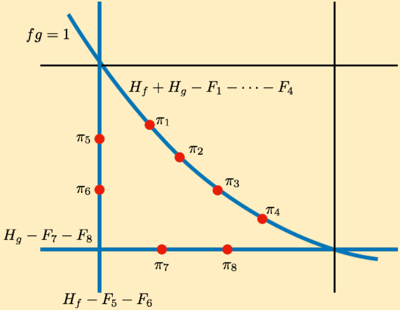

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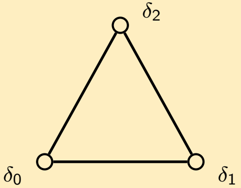

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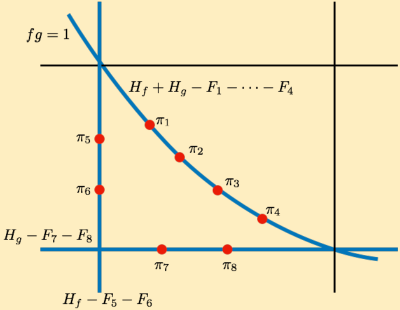

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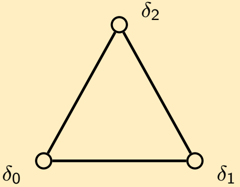

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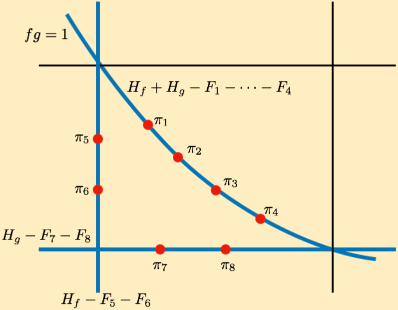

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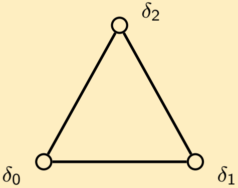

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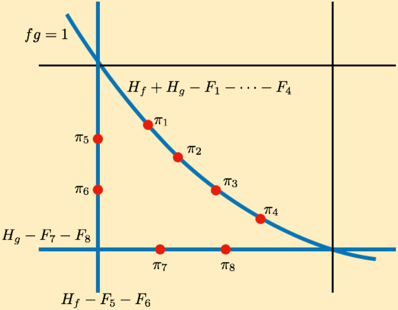

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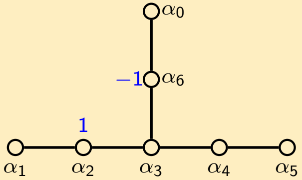
The points π_i lie on the $(2,2)$-curve that is the pole divisor of the symplectic form $\omega = \dfrac{df \wedge dg}{fg(1-fg)} = \dfrac{df \wedge ds}{fs(1-s)} = \dfrac{ds \wedge dg}{gs(1-s)}$, $s = fg$, that is used to define the *period map*.

$E_6^{(1)}$ symmetry sub-lattice $Q = \mathrm{Span}_{\mathbb{Z}}\{\alpha_i | \alpha_i \bullet \delta_j = 0\}$



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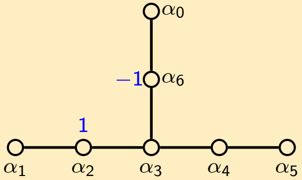

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$A_2^{(1)}/E_6^{(1)}$ period map

The period map $\chi : Q \to \mathbb{C}$, $\chi(\alpha_i) = a_i$, is used to pass from the original parameters ν_i and k_j that still have some Möbius gauge freedom to the invariant *root variables* $a_i = \exp(\alpha_i)$. Moreover, the evolution of the root variables is also canonical. We get

$$a_0 = \frac{\nu_7}{\nu_8}, \quad a_1 = \frac{\nu_6}{\nu_5}, \quad a_2 = \frac{k_2}{\nu_1\nu_6}, \quad a_3 = \frac{\nu_1}{\nu_2}, \quad a_4 = \frac{\nu_2}{\nu_3}, \quad a_5 = \frac{\nu_3}{\nu_4}, \quad a_6 = \frac{k_1}{\nu_1\nu_7}.$$

The dynamic on parameters $\bar{\nu}_i = \nu_i$, $\bar{k}_1 = q^{-1}k_1$, $\bar{k}_2 = qk_2$ results in $\bar{a}_2 = qa_2$, $\bar{a}_6 = q^{-1}a_6$, and $\bar{a}_i = a_i$ otherwise; here $q = \exp(\chi(\delta)) = a_0a_1a_2^2a_3^3a_4^2a_5a_6^2 = \dfrac{k_1k_2}{\nu_1\cdots\nu_8}$.

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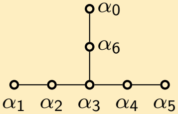
$$W(E_6^{(1)}) = \left\langle w_0, \ldots, w_6 \;\middle|\; \begin{array}{ll} w_i^2 = e & \\ w_i \circ w_j = w_j \circ w_i & \text{when } \alpha_i \quad \alpha_j \\ w_i \circ w_j \circ w_i = w_j \circ w_i \circ w_j & \text{when } \alpha_i - \alpha_j \end{array} \right\rangle$$



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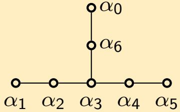

- The finite group of Dynkin diagram automorphisms

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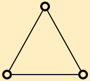

where $\mathbb{D}_3 = \{e, m_0, m_1, m_2, r, r^2\} = \langle m_0, r \mid m_0^2 = r^3 = e, m_0 r = r^2 m_0 \rangle$ is the usual *dihedral group* of the symmetries of a triangle.

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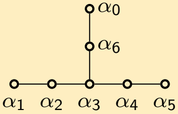

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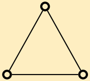

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The action of $\widetilde{W}\left(E_6^{(1)}\right)$ on $\mathrm{Pic}(\mathcal{X})$ can be extended to the action on the space of initial conditions, giving us the birational representation of $\widetilde{W}\left(E_6^{(1)}\right)$.

For the standard example, knowing the action on the *root variables*, $\bar{a}_2 = qa_2$, $\bar{a}_6 = q^{-1}a_6$, and $\bar{a}_i = a_i$ otherwise, we see that mapping φ_* induces the translation

$$\langle \bar{\alpha}_0, \bar{\alpha}_1, \bar{\alpha}_2, \bar{\alpha}_3, \bar{\alpha}_4, \bar{\alpha}_5, \bar{\alpha}_6 \rangle = \langle \alpha_0, \alpha_1, \alpha_2, \alpha_3, \alpha_4, \alpha_5, \alpha_6 \rangle + \langle 0, 0, 1, 0, 0, 0, -1 \rangle \delta$$

and then, using some standard techniques, we can represent this translation as a word in the generators:

$$\varphi_* = r w_2 w_3\, w_1 w_2 w_6\, w_3 w_4 w_0\, w_6 w_3 w_5\, w_4 w_2 w_3\, w_1 w_2.$$

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$$\left(q\text{-P}\left(A_2^1/E_6^{(1)}\right)\right): \quad \begin{cases} \dfrac{(fg-1)(\bar{f}g-1)}{f\bar{f}} = \dfrac{\left(g-\frac{1}{\nu_1}\right)\left(g-\frac{1}{\nu_2}\right)\left(g-\frac{1}{\nu_3}\right)\left(g-\frac{1}{\nu_4}\right)}{\left(g-\frac{\nu_5}{k_2}\right)\left(g-\frac{\nu_6}{k_2}\right)} \\[2ex] \dfrac{(fg-1)(f\underline{g}-1)}{g\underline{g}} = \dfrac{(f-\nu_1)(f-\nu_2)(f-\nu_3)(f-\nu_4)}{\left(f-\frac{k_1}{\nu_7}\right)\left(f-\frac{k_1}{\nu_8}\right)} \end{cases}.$$

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$$\left(q\text{-P}\left(A_2^1/E_6^{(1)}\right)\right): \quad \begin{cases} \dfrac{(fg-1)(\bar{f}g-1)}{f\bar{f}} = \dfrac{\left(g-\frac{1}{\nu_1}\right)\left(g-\frac{1}{\nu_2}\right)\left(g-\frac{1}{\nu_3}\right)\left(g-\frac{1}{\nu_4}\right)}{\left(g-\frac{\nu_5}{k_2}\right)\left(g-\frac{\nu_6}{k_2}\right)} \\[2ex] \dfrac{(fg-1)(f\underline{g}-1)}{g\underline{g}} = \dfrac{(f-\nu_1)(f-\nu_2)(f-\nu_3)(f-\nu_4)}{\left(f-\frac{k_1}{\nu_7}\right)\left(f-\frac{k_1}{\nu_8}\right)} \end{cases}.$$

Note that a more traditional approach is to start with the equation and then obtain the corresponding translation vector.

The q-Hahn Connections and Modui Space Parameterization

Structure of a generic $A(z)$ of type $\lambda = (z_1, \ldots, z_6; u, v, w, w; 3)$

$$A(z) = \begin{bmatrix} a_{11}(z) & a_{12}(z) \\ a_{21}(z) & a_{22}(z) \end{bmatrix}, \quad A(0) = \begin{bmatrix} w & 0 \\ 0 & w \end{bmatrix},$$

where $\deg(a_{11}) \leq 3$, $\deg(a_{12}) \leq 2$, $\deg(a_{21}) \leq 2$, $\deg(a_{22}) \leq 3$ and

$$\det A(z) = uv(z - z_1)(z - z_2)(z - z_3)(z - z_4)(z - z_5)(z - z_6)$$

We also impose asymptotic conditions

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Parameter evolution

When $A(z) \to \overline{A}(z)$, the parameters evolve as

$$(z_1, z_2, \ldots, z_6, u_s, v_s, w_s) \to (\overline{z}_1, \overline{z}_2, \ldots, \overline{z}_6, \overline{u}, \overline{v}, \overline{w})$$

with $\overline{z}_2 = qz_2$, $\overline{z}_4 = qz_4$, $\overline{w} = qw$, and $\overline{z}_i = z_i$ for $i \neq 2, 4$.

Let us now explicitly describe the moduli space of q-Hahn connections of type $\lambda = (z_1, \ldots, z_6; u, qv, w, w, ; 3)$. After gauging we can put $a_{21}(z) = z(z - t)$, where $t = t_1/t_2$ is our first spectral coordinate. The second spectral coordinate we adjust slightly and put

$$p = \frac{p_1}{p_2} = \frac{z_1 z_3 z_5\, a_{11}(t)}{(t - z_1)(t - z_3)(t - z_5)}.$$

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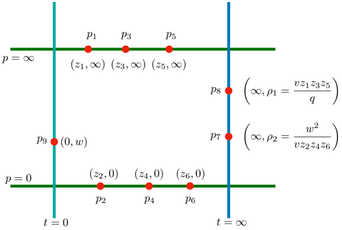
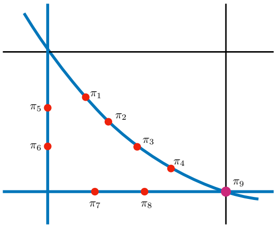


$\xrightarrow{\hspace{4em}}$ need to blow-up at π_9



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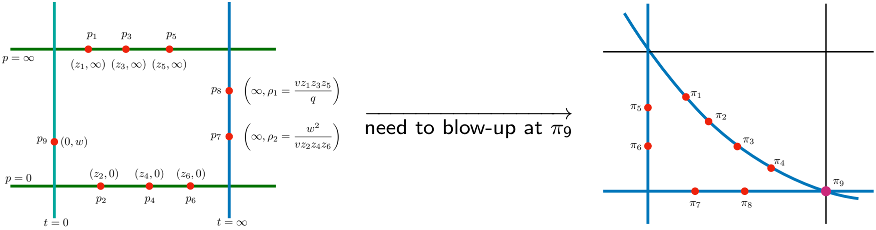

Note that the q-Hahn surface is *not minimal* and requires *blowing down* the -1-curve $t = 0$. It is easier to match it with the standard example by blowing up the point $\pi_9(\infty, 0)$ in the standard (f, g)-coordinates and establishing the identification on the level of Picard lattices, and then extending it to the birational change of coordinates.

Matching the two dynamics

After some minor trial and error, we see that the following identification works:

$$\begin{aligned}
&\mathcal{H}_f = \mathcal{H}_t && \mathcal{F}_1 = \mathcal{E}_1, && \mathcal{F}_3 = \mathcal{E}_3, && \mathcal{F}_5 = \mathcal{E}_7, && \mathcal{F}_7 = \mathcal{E}_2, && \mathcal{F}_9 = \mathcal{H}_t - \mathcal{E}_9, \\
&\mathcal{H}_g = \mathcal{H}_t + \mathcal{H}_p - \mathcal{E}_6 - \mathcal{E}_9, && \mathcal{F}_2 = \mathcal{H}_t - \mathcal{E}_6, && \mathcal{F}_4 = \mathcal{E}_5, && \mathcal{F}_6 = \mathcal{E}_8, && \mathcal{F}_8 = \mathcal{E}_4.
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The standard techniques then give us the explicit change of variables from the *application coordinates* (or the spectral coordinates t and p) to the *Painlevé coordinates* f and g:

$$f = \frac{1}{t}, \qquad g = \frac{twz_6}{z_6(p-w)+tw}.$$

We also get the parameter matching;

$$\begin{aligned}
&k_1 = \frac{1}{w}, & &\nu_1 = \frac{1}{z_1}, & &\nu_3 = \frac{1}{z_3}, & &\nu_5 = \rho_1 z_6, & &\nu_7 = \frac{z_2}{w},\\
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\end{aligned}$$

(note that there is a parameter constraint in q-Hahn, $w^2 = uvz_1 \cdots z_6$).

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(note that there is a parameter constraint in q-Hahn, $w^2 = uvz_1\cdots z_6$).
With this identification the spectral coordinates evolution under isomonodromic transformations coincides with q-P $\left(A_2^1/E_6^{(1)}\right)$ of (**KNY**).

The q-Racah orthogonal ensemble and q-$P(A_1^{(1)}/E_7^{(1)})$

Consider now the example of a q-Racah orthogonal polynomial ensemble and q-$P\left(A_1^{(1)}/E_7^{(1)}\right)$.

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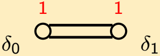
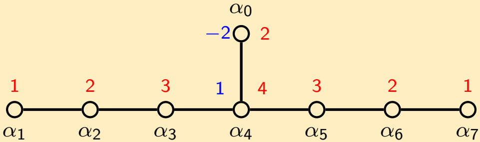
q-$P\left(A_1^{(1)}/E_7^{(1)}\right)$ surface and reference dynamic



Dynkin diagram $A_1^{(1)}$

$$\delta_0 = \mathcal{H}_f + \mathcal{H}_g - \mathcal{F}_1 - \mathcal{F}_2 - \mathcal{F}_3 - \mathcal{F}_4$$
$$\delta_1 = \mathcal{H}_f + \mathcal{H}_g - \mathcal{F}_5 - \mathcal{F}_6 - \mathcal{F}_7 - \mathcal{F}_8$$

The surface data



Dynkin diagram $E_7^{(1)}$

$$\alpha_0 = \mathcal{H}_f - \mathcal{H}_g \qquad \alpha_4 = \mathcal{H}_g - \mathcal{F}_1 - \mathcal{F}_5$$
$$\alpha_1 = \mathcal{F}_3 - \mathcal{F}_4 \qquad \alpha_5 = \mathcal{F}_5 - \mathcal{F}_6$$
$$\alpha_2 = \mathcal{F}_2 - \mathcal{F}_3 \qquad \alpha_6 = \mathcal{F}_6 - \mathcal{F}_7$$
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The symmetry data

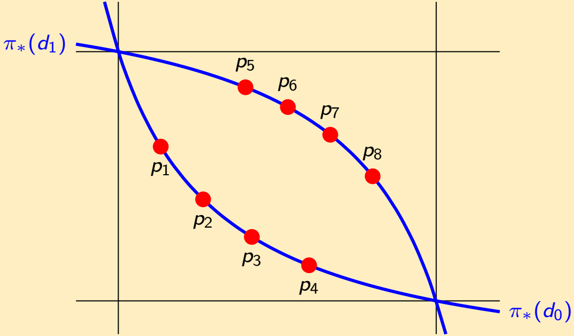
$A_1^{(1)}$ point configuration and the surface model

$$p_i\left(\nu_i, \frac{1}{\nu_i}\right), \quad i = 1, \ldots, 4;$$

$$p_i\left(\frac{\kappa_1}{\nu_i}, \frac{\nu_i}{\kappa_2}\right), \quad i = 5, \ldots, 8.$$

$$\pi_*(d_0):\ fg = 1$$

$$\pi_*(d_1):\ fg = \kappa = \frac{\kappa_1}{\kappa_2}$$



The points π_i lie on the (reducible) $(2,2)$-curve that is the pole divisor of the symplectic form $\omega = (k-1)\dfrac{df \wedge dg}{(fg-1)(fg-k)} = (k-1)\dfrac{df \wedge ds}{f(s-1)(s-k)} = (k-1)\dfrac{ds \wedge dg}{g(s-1)(s-k)}$, where again we put $s = fg$. Degeneration to $\left(A_1^{(1)}/E_7^{(1)}\right)$ case is very straightforward, just put $\kappa \to 0$.

$A_1^{(1)}/E_1^{(1)}$ period map

The period map $\chi : Q \to \mathbb{C}$, $\chi(\alpha_i) = a_i$, in this case gives us the root variables $a_i = \exp(\chi(\alpha_i))$:

$$a_0 = \frac{\kappa_1}{\kappa_2}, \quad a_1 = \frac{\nu_3}{\nu_4}, \quad a_2 = \frac{\nu_2}{\nu_3}, \quad a_3 = \frac{\nu_1}{\nu_2}, \quad a_4 = \frac{\kappa_2}{\nu_1 \nu_5}, \quad a_5 = \frac{\nu_5}{\nu_6}, \quad a_6 = \frac{\nu_6}{\nu_7}, \quad a_7 = \frac{\nu_7}{\nu_8}.$$

The dynamic on parameters $\bar{\nu}_i = \nu_i$, $\bar{k}_1 = q^{-1} k_1$, $\bar{k}_2 = q k_2$ results in $\bar{a}_0 = q^{-2} a_0$, $\bar{a}_4 = q a_4$, and $\bar{a}_i = a_i$ otherwise.

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For the standard example, we can represent the mapping φ_* that induces the translation

$$\varphi_* : \alpha = \langle \alpha_0, \alpha_1, \alpha_2, \alpha_3, \alpha_4, \alpha_5, \alpha_6, \alpha_7 \rangle \mapsto \overline{\alpha} = \alpha + \langle 2, 0, 0, 0, -1, 0, 0, 0 \rangle \delta.$$

as

$$\varphi_* : w_0 w_4 w_5 w_3 w_4 w_6 w_5 w_2 w_3 w_4 w_1 w_2 w_3 w_0 w_4 w_7 w_6 w_5 w_4 w_3 w_0 w_4 w_6 w_5 w_2 w_3 w_4 w_7 w_6 w_5 w_1 w_2 w_3 w_4,$$

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$$\begin{cases} \dfrac{\left(fg - \frac{\kappa_1}{\kappa_2}\right)\left(\overline{f}g - \frac{\kappa_1}{q\kappa_2}\right)}{(fg-1)(\overline{f}g-1)} = \dfrac{\left(g - \frac{\nu_5}{\kappa_2}\right)\left(g - \frac{\nu_6}{\kappa_2}\right)\left(g - \frac{\nu_7}{\kappa_2}\right)\left(g - \frac{\nu_8}{\kappa_2}\right)}{\left(g - \frac{1}{\nu_1}\right)\left(g - \frac{1}{\nu_2}\right)\left(g - \frac{1}{\nu_3}\right)\left(g - \frac{1}{\nu_4}\right)}, \\[2ex] \dfrac{\left(fg - \frac{\kappa_1}{\kappa_2}\right)\left(f\underline{g} - \frac{q\kappa_1}{\kappa_2}\right)}{(fg-1)(f\underline{g}-1)} = \dfrac{\left(f - \frac{\kappa_1}{\nu_5}\right)\left(f - \frac{\kappa_1}{\nu_6}\right)\left(f - \frac{\kappa_1}{\nu_7}\right)\left(f - \frac{\kappa_1}{\nu_8}\right)}{(f-\nu_1)(f-\nu_2)(f-\nu_3)(f-\nu_4)}. \end{cases}$$

Moduli space for q-Racah connections

In this case, we look at moduli spaces $\lambda = (z_1, \ldots, z_6; u, d = d_1, d_2; 6)$ of 2×2 matrices satisfying the following conditions:

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with some asymptotic conditions and modulo gauge transformations of the form

$$\hat{A}(z) = R(z/q + u^2/z)A(z)R^{-1}(z + u^2/(qz)), \qquad R(z) = \begin{bmatrix} r_{11}(z) & r_{12}(z) \\ 0 & r_{22}(z) \end{bmatrix},$$

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$$\bar{z}_2 = qz_2, \quad \bar{z}_4 = qz_4, \quad \bar{d} = q^{-1}d, \quad \bar{z}_i = z_i \text{ otherwise.}$$

Moduli space of q-Racah connections

Let us now explicitly describe the moduli space of q-Racah connections. After gauging we can put $a_{21}(z) = z(z-t)(z-u^2)(z^2-u^2)$, where $t = t_1/t_2$ is our first spectral coordinate, and the second spectral coordinate p is again the *adjusted* value of $a_{11}(t)$.

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In the coordinates (t,p) we get more than 8 points because of the involution $t \leftrightarrow u^2/t$ and $p \leftrightarrow 1/p$, e.g., we get the following six pairs of points:

$$\left(\frac{u^2}{z_1},0\right),(z_1,\infty), \quad \left(\frac{u^2}{z_3},0\right),(z_3,\infty), \quad \left(\frac{u^2}{z_5},0\right),(z_5,\infty),$$
$$(z_2,0),\left(\frac{u^2}{z_2},\infty\right), \quad (z_4,0),\left(\frac{u^2}{z_4},\infty\right), \quad (z_6,0),\left(\frac{u^2}{z_6},\infty\right),$$

points $(u,1)$ and $(-u,-1)$, and points $(\infty,-\rho_1=d)$ and $\left(\infty,-\rho_2=\dfrac{z_1z_3z_5}{z_2z_4z_6qd}\right)$.

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To fix this, we need to introduce the *involution-invariant* coordinates $x = t+\dfrac{u^2}{t}$ and $y = \dfrac{pt-u}{pu-t}$ gluing these pairs of points together.

Then, in the (x,y)-coordinates we get the correct picture.

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Then, in the (x,y)-coordinates we get the correct picture.

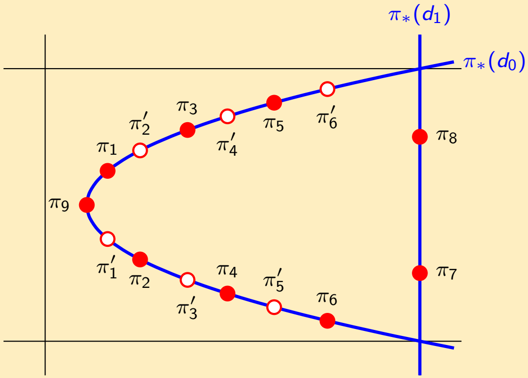
But can it be matched with the standard example?

The q-Racah surface

$$\pi_i\left(z_i+\frac{u^2}{z_i},\frac{z_i}{u}\right),\quad i=1,3,5;$$
$$\pi_i\left(z_i+\frac{u^2}{z_i},\frac{u}{z_i}\right),\quad i=2,4,6;$$
$$\pi_7\left(\infty,\rho_1=-d\right),$$
$$\pi_8\left(\infty,\rho_2=-\frac{z_1z_3z_5}{z_2z_4z_6qd}\right).$$

We also get conjugated points

$$\pi_i'\left(z_i+\frac{u^2}{z_i},\frac{u}{z_i}\right),\quad i=1,3,5;$$
$$\pi_i'\left(z_i+\frac{u^2}{z_i},\frac{z_i}{u}\right),\quad i=2,4,6.$$



Note that the points π_7 and π_8 lie on the $(1,0)$-curve $\pi_*(d_1)=V(X=1/x)$ and $\pi_1,\ldots,\pi_6$ lie on the $(1,2)$-curve $\pi_*(d_0)=V(u(y^2+1)-xy)$; note also that when $x=z_i+\dfrac{u^2}{z_i}$, the equation $u(y^2+1)-xy$ factors as $u(y^2+1)-xy=u(y-y(\pi_i))(y-y(\pi_i'))$. Finally, there is an additional blowup point $\pi_9(-2u,-1)$, similar to the q-Hahn case.

Matching the q-Racah dynamics with the standard dynamics

Looking at the decomposition of the anti-canonical divisor class,

$$\delta_0 = \mathcal{H}_f + \mathcal{H}_g - \mathcal{F}_1 - \mathcal{F}_2 - \mathcal{F}_3 - \mathcal{F}_4 - \mathcal{F}_9 = \mathcal{H}_x + 2\mathcal{H}_y - \mathcal{E}_1 - \mathcal{E}_2 - \mathcal{E}_3 - \mathcal{E}_4 - \mathcal{E}_5 - \mathcal{E}_6 - \mathcal{E}_9,$$
$$\delta_1 = \mathcal{H}_f + \mathcal{H}_g - \mathcal{F}_5 - \mathcal{F}_6 - \mathcal{F}_7 - \mathcal{F}_8 = \mathcal{H}_x - \mathcal{E}_7 - \mathcal{E}_8,$$

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\delta_1 &= \mathcal{H}_f + \mathcal{H}_g - \mathcal{F}_5 - \mathcal{F}_6 - \mathcal{F}_7 - \mathcal{F}_8 = \mathcal{H}_x - \mathcal{E}_7 - \mathcal{E}_8,
\end{aligned}$$

we see that it makes sense to preliminary take

$$\begin{aligned}
\mathcal{H}_f &= \mathcal{H}_x + \mathcal{H}_y - \mathcal{E}_2 - \mathcal{E}_9, & \mathcal{H}_x &= \mathcal{H}_f + \mathcal{H}_g - \mathcal{F}_7 - \mathcal{F}_8,\\
\mathcal{H}_g &= \mathcal{H}_x + \mathcal{H}_y - \mathcal{E}_4 - \mathcal{E}_9, & \mathcal{H}_y &= \mathcal{H}_f + \mathcal{H}_g - \mathcal{F}_7 - \mathcal{F}_9,\\
\mathcal{F}_1 &= \mathcal{E}_1, & \mathcal{E}_1 &= \mathcal{F}_2,\\
\mathcal{F}_2 &= \mathcal{E}_6, & \mathcal{E}_2 &= \mathcal{H}_g - \mathcal{F}_7,\\
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\mathcal{F}_4 &= \mathcal{E}_5, & \mathcal{E}_4 &= \mathcal{H}_f - \mathcal{F}_7,\\
\mathcal{F}_5 &= \mathcal{E}_7, & \mathcal{E}_5 &= \mathcal{F}_4,\\
\mathcal{F}_6 &= \mathcal{E}_8, & \mathcal{E}_6 &= \mathcal{F}_1,\\
\mathcal{F}_7 &= \mathcal{H}_x + \mathcal{H}_y - \mathcal{E}_2 - \mathcal{E}_4 - \mathcal{E}_9, & \mathcal{E}_7 &= \mathcal{F}_5,\\
\mathcal{F}_8 &= \mathcal{H}_y - \mathcal{E}_9, & \mathcal{E}_8 &= \mathcal{F}_6,\\
\mathcal{F}_9 &= \mathcal{H}_x - \mathcal{E}_9, & \mathcal{E}_9 &= \mathcal{H}_f + \mathcal{H}_g - \mathcal{F}_7 - \mathcal{F}_8 - \mathcal{F}_9.
\end{aligned}$$

Matching the q-Racah dynamics with the standard dynamics

Computing the root variables,

$$a_0 = \frac{z_4}{z_2},\ a_1 = \frac{z_5}{z_3},\ a_2 = \frac{z_3}{z_1},\ a_3 = \frac{z_1 z_6}{u^2},\ a_4 = -\frac{u^2}{\rho_1 z_4 z_6},\ a_5 = \frac{\rho_1}{\rho_2},\ a_6 = -\frac{\rho_2 z_2 z_4}{u^2},\ a_7 = \frac{u^2}{z_2 z_4}.$$

and using our parameter dynamics, we get the following translation element:

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$$\psi_* : \alpha = \langle \alpha_0, \alpha_1, \alpha_2, \alpha_3, \alpha_4, \alpha_5, \alpha_6, \alpha_7 \rangle \mapsto \overline{\alpha} = \alpha + \langle 0,0,0,0,0,0,-1,2\rangle\delta,$$

which is **different** from the standard translation vector

$$\varphi_* : \alpha = \langle \alpha_0, \alpha_1, \alpha_2, \alpha_3, \alpha_4, \alpha_5, \alpha_6, \alpha_7 \rangle \mapsto \overline{\alpha} = \alpha + \langle 2,0,0,0,-1,0,0,0\rangle\delta.$$

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However, these elements are **conjugated**. This can be observed, for example, by looking at the corresponding words in the affine Weyl symmetry group:

$$\psi_* : w_7 w_6 w_5 w_4 w_3 w_0 w_4 \color{red}{w_5 w_2 w_3 w_4 w_1 w_2 w_3 w_0 w_4 w_6 w_5 w_4 w_3 w_0 w_4 w_6 w_5 w_2 w_3 w_4 w_1 w_2 w_3} w_0 w_4 w_5 w_6,$$

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$$\varphi_* : w_0 w_4 w_5 w_3 w_4 w_6 w_5 w_2 w_3 w_4 w_1 w_2 w_3 w_0 w_4 w_7 w_6 w_5 w_4 w_3 w_0 w_4 w_6 w_5 w_2 w_3 w_4 w_7 w_6 w_5 w_1 w_2 w_3 w_4.$$

Using the far commutativity and the braid relations in $W\left(E_7^{(1)}\right)$, we get

$$\psi_* = (w_6 w_5 w_4 w_0 w_7 w_6 w_5 w_4) \varphi_* (w_6 w_5 w_4 w_0 w_7 w_6 w_5 w_4)^{-1}.$$

Matching the q-Racah dynamics with the standard dynamics

This adjusts the divisor matching as

$$
\begin{aligned}
\mathcal{H}_f &= 2\mathcal{H}_x + \mathcal{H}_y - \mathcal{E}_2 - \mathcal{E}_4 - \mathcal{E}_6 - \mathcal{E}_9, & \mathcal{H}_x &= \mathcal{H}_f + \mathcal{H}_g - \mathcal{F}_5 - \mathcal{F}_6,\\
\mathcal{H}_g &= \mathcal{H}_x + \mathcal{H}_y - \mathcal{E}_6 - \mathcal{E}_9, & \mathcal{H}_y &= \mathcal{H}_f + 2\mathcal{H}_g - \mathcal{F}_1 - \mathcal{F}_5 - \mathcal{F}_6 - \mathcal{F}_9,\\
\mathcal{F}_1 &= \mathcal{H}_x - \mathcal{E}_6, & \mathcal{E}_1 &= \mathcal{F}_2,\\
\mathcal{F}_2 &= \mathcal{E}_1, & \mathcal{E}_2 &= \mathcal{H}_g - \mathcal{F}_5,\\
\mathcal{F}_3 &= \mathcal{E}_3, & \mathcal{E}_3 &= \mathcal{F}_3,\\
\mathcal{F}_4 &= \mathcal{E}_5, & \mathcal{E}_4 &= \mathcal{H}_g - \mathcal{F}_6,\\
\mathcal{F}_5 &= \mathcal{H}_x + \mathcal{H}_y - \mathcal{E}_2 - \mathcal{E}_6 - \mathcal{E}_9, & \mathcal{E}_5 &= \mathcal{F}_4,\\
\mathcal{F}_6 &= \mathcal{H}_x + \mathcal{H}_y - \mathcal{E}_4 - \mathcal{E}_6 - \mathcal{E}_9, & \mathcal{E}_6 &= \mathcal{H}_f + \mathcal{H}_g - \mathcal{F}_1 - \mathcal{F}_5 - \mathcal{F}_6,\\
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\mathcal{F}_9 &= \mathcal{H}_x - \mathcal{E}_9, & \mathcal{E}_9 &= \mathcal{H}_f + \mathcal{H}_g - \mathcal{F}_5 - \mathcal{F}_6 - \mathcal{F}_9.
\end{aligned}
$$

Matching the q-Racah dynamics with the standard dynamics

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$$
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\mathcal{H}_f &= 2\mathcal{H}_x + \mathcal{H}_y - \mathcal{E}_2 - \mathcal{E}_4 - \mathcal{E}_6 - \mathcal{E}_9, & \mathcal{H}_x &= \mathcal{H}_f + \mathcal{H}_g - \mathcal{F}_5 - \mathcal{F}_6,\\
\mathcal{H}_g &= \mathcal{H}_x + \mathcal{H}_y - \mathcal{E}_6 - \mathcal{E}_9, & \mathcal{H}_y &= \mathcal{H}_f + 2\mathcal{H}_g - \mathcal{F}_1 - \mathcal{F}_5 - \mathcal{F}_6 - \mathcal{F}_9,\\
\mathcal{F}_1 &= \mathcal{H}_x - \mathcal{E}_6, & \mathcal{E}_1 &= \mathcal{F}_2,\\
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\mathcal{F}_5 &= \mathcal{H}_x + \mathcal{H}_y - \mathcal{E}_2 - \mathcal{E}_6 - \mathcal{E}_9, & \mathcal{E}_5 &= \mathcal{F}_4,\\
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\mathcal{F}_9 &= \mathcal{H}_x - \mathcal{E}_9, & \mathcal{E}_9 &= \mathcal{H}_f + \mathcal{H}_g - \mathcal{F}_5 - \mathcal{F}_6 - \mathcal{F}_9.
\end{aligned}
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The new root variables are

$$a_0 = \frac{u^2}{z_2 z_4},\ a_1 = \frac{z_5}{z_3},\ a_2 = \frac{z_3}{z_1},\ a_3 = \frac{z_1}{z_6},\ a_4 = \frac{z_2 z_6}{u^2},\ a_5 = \frac{z_4}{z_2},\ a_6 = -\frac{u^2}{\rho_1 z_4 z_6},\ a_7 = \frac{\rho_1}{\rho_2},$$

which, given the parameter dynamic $\bar{z}_2 = q z_2$, $\bar{z}_4 = q z_4$, $\bar{d} = q^{-1} d$ (and so $\bar{\rho}_i = q^{-1}\rho_i$), immediately gives us the correct translation element:

$$\psi_* : \alpha = \langle \alpha_0, \alpha_1, \alpha_2, \alpha_3, \alpha_4, \alpha_5, \alpha_6, \alpha_7 \rangle \mapsto \overline{\alpha} = \alpha + \langle 2, 0, 0, 0, -1, 0, 0, 0 \rangle \delta.$$

Matching the q-Racah dynamics with the standard dynamics

We then get the following parameter matching: $\kappa_1 = \dfrac{u}{z_2}$, $\kappa_2 = \dfrac{z_4}{u}$, and

$$\nu_1 = \frac{1}{z_6},\ \nu_2 = \frac{1}{z_1},\ \nu_3 = \frac{1}{z_3},\ \nu_4 = \frac{1}{z_5},\ \nu_5 = \frac{u z_4}{z_2},\ \nu_6 = u,\ \nu_7 = -\frac{\rho_1 z_4 z_6}{u},\ \nu_8 = -\frac{\rho_2 z_4 z_6}{u}.$$

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Main Result

The change of variables from the spectral coordinates to the discrete Painlevé coordinates matching the q-Racah isomonodromic dynamics to the standard dynamics is given by

$$\begin{aligned} f(x,y) &= \frac{\sigma_3(xy + u(y-1)) - u^2(x^2 - \sigma_1 x + \sigma_2(y+1)) + u^3(1-y)(\sigma_1 - x) + u^4(1+y)}{\sigma_3 x(xy + u(y-1)) - u^2(\sigma_2 xy + \sigma_3(y+1)) + u^3\sigma_2(1-y) + u^4(\sigma_1(1+y) - x) + u^5(y-1)}, \\ &\to \frac{1}{x}, \\ g(x,y) &= \frac{xyz_6 + uz_6(y-1) - u^2(1+y)}{z_6(1+y) - x - u(1+y)} \to \frac{xyz_6}{z_6(1+y) - x}, \qquad \text{where} \\ &\sigma_1 = z_2 + z_4 + z_6, \qquad \sigma_2 = z_2 z_4 + z_4 z_6 + z_6 z_2, \qquad \sigma_3 = z_2 z_4 z_6. \end{aligned}$$

Everything does have the correct limit to the q-Hahn case as $u \to 0$.

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