# Quantum Walking in Curved Spacetime 

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Joint work with S. Facchini, M. Forets.

## What's in the tin?

A stable numerical scheme for PDEs of the form

$$
i \partial_{0} \psi=H \psi
$$

$$
H=i \sum_{k=1 \ldots d}\left(B_{k} \partial_{k}+\frac{1}{2} \partial_{k} B_{k}\right)-C
$$

$$
\left(\text { with } B_{k}, C \in \operatorname{Herm}(\mathbb{C}) \text { and }\left|B_{k}\right| \leq 1\right)
$$

implementable by applying unitary matrices locally i.e. by future quantum simulation devices.

## Discretize physics?



Cellular Automata
An old CompSci dream : to capture physics in this formalism.

## Discretize physics?


... as Cellular Automata / Quantum Walks
Theorems about : the extent in which physics particles can be captured in this formalism.

## Discretize particules

Dirac equation
$\mathrm{i} \partial_{0} \psi=D \psi, \quad$ with $\quad D=m \alpha^{0}-\mathrm{i} \sum_{j} \alpha^{j} \partial_{j}$
vs

Chess game


## Chess game : neutrino

To the right

| \& | $\binom{0}{0}$ | $\binom{0}{0}$ | $\binom{0}{0}$ | $\binom{0}{0}$ |
| :---: | :---: | :---: | :---: | :---: |
| $\binom{0}{0}$ | $\binom{0}{0}$ | $\binom{0}{0}$ | $\binom{0}{0}$ | $\binom{1}{0}$ |
| $\binom{0}{0}$ | $\binom{0}{0}$ | $\binom{0}{0}$ | 1 | $\binom{0}{0}$ |
| $\binom{0}{0}$ | $\binom{0}{0}$ | $\left(\begin{array}{l}1 \\ 0\end{array}\right.$ | $\binom{0}{0}$ | $\binom{0}{0}$ |
| $\binom{0}{0}$ | 1 | $\binom{0}{0}$ | $\binom{0}{0}$ | $\binom{0}{0}$ |

## Chess game : neutrino

To the right

| $\stackrel{\otimes}{\otimes}\binom{0}{0}$ | $\binom{0}{0}$ | $\binom{0}{0}$ | $\binom{0}{0}$ | $\binom{0}{0}$ |
| :---: | :---: | :---: | :---: | :---: |
| 0 | $\binom{0}{0}$ | $\binom{0}{0}$ | $\binom{0}{0}$ | $\binom{0}{0}$ |
| $\binom{0}{0}$ |  | $\binom{0}{0}$ | $\binom{0}{0}$ | $\binom{0}{0}$ |
| $\binom{0}{0}$ | $\binom{0}{0}$ |  | $\binom{0}{0}$ | $\binom{0}{0}$ |
| $\binom{0}{0}$ | $\binom{0}{0}$ | $\binom{0}{0}$ |  | $\binom{0}{0}$ |

## Chess game : neutrino

To the right

To the left

Amplitudes

$$
|\alpha|^{2}+|\beta|^{2}=1
$$



| $\stackrel{\otimes}{ \pm}\binom{0}{0}$ | $\binom{0}{0}$ | $\binom{0}{0}$ | $\binom{0}{0}$ | $\binom{0}{0}$ |
| :---: | :---: | :---: | :---: | :---: |
| $\binom{0}{\beta}$ | $\binom{0}{0}$ | $\binom{0}{0}$ | $\binom{0}{0}$ | $\binom{a}{0}$ |
| $\binom{0}{0}$ | $\binom{0}{\beta}$ | $\binom{0}{0}$ | $\binom{a}{0}$ | $\binom{0}{0}$ |
| $\binom{0}{0}$ | $\binom{0}{0}$ | $\binom{a}{\beta}$ | $\binom{0}{0}$ | $\binom{0}{0}$ |
| $\binom{0}{0}$ | $\binom{a}{0}$ | $\binom{0}{0}$ | $\binom{0}{\beta}$ | $\binom{0}{0}$ |

## Chess game : electron

Rotations
$C=\left(\begin{array}{l}C \\ s\end{array}\right.$
$c=\cos (\theta)$
$s=\sin (\theta)$
$\theta=m . \varepsilon$
$m=$ mass
$\varepsilon=$ step

| $\stackrel{\mathbb{1}}{\underset{=}{E}}$ | $\binom{0}{0}$ | $\binom{\ldots}{\ldots}$ | $\binom{0}{0}$ | $\left(\begin{array}{l}. \\ \ldots \\ \ldots\end{array}\right)$ | $\binom{0}{0}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\binom{-c s^{2}}{c^{2} s}$ | $\binom{0}{0}$ | $\binom{-2 c s^{2}}{-s^{3}+c^{2} s}$ | $\binom{0}{0}$ | $\binom{c^{3}}{c^{2} s}$ |
|  | $\binom{0}{0}$ | $\binom{-s^{2}}{c s}$ | $\binom{0}{0}$ | $\binom{c^{2}}{c s}$ | $\binom{0}{0}$ |
|  | $\binom{0}{0}$ | $\binom{0}{0}$ | $\binom{c}{s}$ | $\binom{0}{0}$ | $\binom{0}{0}$ |
|  | $\binom{0}{0}$ | $\binom{1}{0}$ | $\binom{0}{0}$ | $\binom{0}{0}$ | $\binom{0}{0}$ |

## Chess game : electron

Rotations


## Chess game : electron

States

Scale


## Chess game : electron



## Chess game : electron

Expand


## Chess game : electron

Order 0


## Chess game : electron

Order 0


## Chess game : electron

Order 1

$$
d+\varepsilon \partial_{t} d=d+\varepsilon \partial_{x} d+\ldots \quad u+\varepsilon \partial_{t} u=u-\varepsilon \partial_{x} u+\ldots
$$

## Chess game : electron

Order 1


## Chess game : electron

Order 1


## Chess game : electron

Order 1

$$
\begin{aligned}
& C=\left(\begin{array}{ll}
c & -s \\
s & c
\end{array}\right) \quad \partial_{t}\binom{u}{d}=-\sigma_{z} \partial_{x}\binom{u}{d}+W^{(1)}\binom{u}{d} \\
& \mathrm{c}=\cos (\theta) \\
& \mathrm{s}=\sin (\theta) \\
& \theta=m \cdot \varepsilon \\
& \begin{array}{l}
C^{(1)}=-i m \sigma_{y} \\
W=C X
\end{array}
\end{aligned}
$$

## Chess game : electron

Rotations


## Chess game : electron

Rotations


## Consistency vs convergence



## Consistency vs convergence



## Consistency vs convergence

Theorem

$$
\begin{aligned}
& \forall \psi(0) \in H^{2}, \forall t, \forall \varepsilon: \\
& \quad\left\|W_{\varepsilon}^{\lfloor t / \varepsilon\rfloor} \psi(0)-\psi(t)\right\|_{2}=\varepsilon(5 / 2) t\|\psi(0)\|_{H^{2}}
\end{aligned}
$$

Morale: unitarity gives you stability in Sobolev norm, for free, and so convergence is for free, too.

## Discretize physics?



Cellular Automata / Quantum Walks
Theorems about : the extent in which Curved Spacetime can be captured in this formalism.

## Curved space : problem 1

From


## Curved space : problem 1

To


## Curved space : problem 1



Transport term is fixed by $0^{\text {th }}$ order \& grid :-((

## Curved space : idea 0

[Di Molfetta , F. Debbasch, M. E. Brachet, "Quantum walks as massless Dirac Fermions in curved Space-Time", PRA, arXiv:1212.5821]


## Curved space : idea 0

[Di Molfetta , F. Debbasch, M. E. Brachet, "Quantum walks as massless Dirac Fermions in curved Space-Time", PRA, arXiv:1212.5821]

$$
\partial_{t}\binom{\mathrm{u}}{\mathrm{~d}}=\left(-\sigma_{\mathrm{z}} \partial_{\partial x}\binom{\mathrm{u}}{\mathrm{~d}}\right.
$$

- $1=r=c$
- massless
- $1+1$

Made it

$$
\left(\begin{array}{ll}
-c(x, t) & 0 \\
0 & c(x, t)
\end{array}\right)
$$

## Curved space : idea 1



## Curved space : idea 1



## Curved space : idea 1

States


## Curved space : idea 1

States


## Curved space : idea 1

States


## Curved space : idea 1



## Curved space : idea 1

Expand


## Curved space : idea 1

$0^{\text {th }}$ order


## Curved space : idea 1

$0^{\text {th }}$ order


## Curved space : problem 2

$1^{\text {st }}$ order


## Curved space : problem 2

$1^{\text {st }}$ order


## Curved space : problem 2

$1^{\text {st }}$ order


## Curved space : idea 2



## Curved space : idea 2



## Curved space : idea $2^{w}$

## Curved space : idea $2^{\text {w }}$



## Tin content

Theorem
A stable numerical scheme for PDEs of the form

$$
\begin{aligned}
& \quad \begin{array}{l}
i \partial_{0} \psi=H \psi \\
\quad H=i \sum_{k=1 \ldots d}\left(B_{k} \partial_{k}+\frac{1}{2} \partial_{k} B_{k}\right)-C \\
\left(\text { with } B_{k}, C \in \operatorname{Herm}(\mathbb{C}) \text { and }\left|B_{k}\right| \leq 1\right)
\end{array}
\end{aligned}
$$

implementable by applying unitary matrices locally.

## Curved space simulations : RW



## Curved space simulations: BH



## Conclusion

Non-interacting physics particles in curved space-time ...as a Quantum Walk.

The point?

- stable numerical scheme
- quantum simulation device compatible
- to simplify, understand, offer toy models.

OK, but what about symmetries?

## Extra 1

## Discretize physics?



Cellular Automatas / Quantum Walks
Theorems about : the extent in which the SR notion of time can be captured in this formalism.

## Time in SR

Observer at rest

$$
\begin{aligned}
& \text { Whas wha } \\
& \text { whas }
\end{aligned}
$$



## Time in SR

Observer at rest


## Time in SR

Observer at rest

Uniform observer


## Time in SR

Observer at rest

Uniform observer

Relativity
Both are right.
who Whan

$$
3.10^{8} \mathrm{~m} \cdot \mathrm{~s}^{-1}
$$

Laws of Physics are the same for uniform observers. Any uniform referential is valid for describing the world

## Time in SR



## Time in SR



## Lorentz transform



## Lorentz transform



## Lorentz transform



## Discrete Lorentz transform



## Discrete Lorentz transform



## Discrete Lorentz transform



## Discrete Lorentz transform



## Discrete Lorentz transform



## Discrete Lorentz transform



## Covariance

Relativity
Laws of Physics are the same for uniform observers.
Any uniform referential is valid for describing the world.

Transform(Dirac Equation) = Dirac Equation
Transform(Physics Law) = Physics Law

A fundamental symmetry of physics.
Can it be discretized?

## Covariance

Relativity
Laws of Physics are the same for uniform observers.
Any uniform referential is valid for describing the world.

Transform(Dirac Equation) = Dirac Equation
Transform(Quantum Walk) = Quantum Walk?

A fundamental symmetry of physics.
Can it be discretized?

## Covariance

Transform(Quantum Walk) = Quantum Walk?


## Discrete covariance



## Discrete covariance

If


## Discrete covariance

If

Then
Theorem :

- The Dirac QW is first-order only discrete-covariant.
- The Clock QCA is discrete-covariant and simulates the Dirac QW.


## Indulging into reductionism



## Indulging into reductionism

We might leave in a "great quantum circuit".

This great quantum circuit would be equivalent to some others... each of which would be a valid representation of our world.

The notion of time would then be relative to this choice of representation, just like in SR.

## Extra 2

## Curved space : idea 2

$0^{\text {th }}$ order


## Curved space : idea 2

$0^{\text {th }}$ order


## Curved space : idea 2

$0^{\text {th }}$ order


## Curved space : idea 2

$0^{\text {th }}$ order


$$
E^{\dagger} W^{(0)} X E=I \oplus U
$$

$\mathrm{W}^{(0)}$, that which governs propagation, is non-trivial, and this still has a continuous limit.

## Curved space : idea 2

$1^{\text {st }}$ order


## Curved space : idea 2



## Curved space : idea 2

$1^{\text {st }}$ order

$$
E^{\dagger} W^{(0)} X E=I \oplus U
$$



## Curved space : idea 2

$1^{\text {st }}$ order

$$
\begin{aligned}
& u=\psi^{+} \\
& d=\psi^{-} \\
& u^{\prime}=2 \varepsilon \partial_{x} \psi^{-} \\
& d^{\prime}=2 \varepsilon \partial_{x} \psi^{-}
\end{aligned}
$$

$$
x_{-2 u^{-}-2 d^{\prime} u^{\prime} d^{\prime}}
$$


$E^{\dagger} W^{(0)} X E=I \oplus U$
$W^{(0)}$, that which governs propagation, is non-trivial, and this still has a continuous limit.

Some constraints for consistency? Yes, but with non-trivial solutions.

## Curved space : Dirac Eq.



